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FLORIDA INTERNATIONAL UNIVERSITY

Miami, Florida

SONOCHEMICAL SYNTHESIS OF ZINC OXIDE NANOSTRUCTURES FOR SENSING AND ENERGY HARVESTING

A dissertation submitted in partial fulfillment of

the requirements for the degree of

DOCTOR OF PHILOSOPHY

in

ELECTRICAL ENGINEERING

by

Phani Kiran Vabbina

2016

To: Interim Dean Ranu Jung College of Engineering and Computing

This dissertation, written by Phani Kiran Vabbina, and entitled Sonochemical Synthesis of Zinc Oxide Nanostructures for Sensing and Energy Harvesting, having been approved in respect to style and intellectual content, is referred to you for judgment.

We have read this dissertation and recommend that it be approved.

Chunlei Wang

Gang Quan

Jean Andrian

Nezih Pala, Major Professor

Date of Defense: July 6, 2016

The dissertation of Phani Kiran Vabbina is approved.

Interim Dean Ranu Jung College of Engineering and Computing

Andrés G. Gil Vice President for Research and Economic Development and Dean of the University Graduate School

Florida International University, 2016

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ABSTRACT OF THE DISSERTATION

SONOCHEMICAL SYNTHESIS OF ZINC OXIDE NANOSTRUCTURES FOR SENSING AND ENERGY HARVESTING

by

Phani Kiran Vabbina

Florida International University, 2016

Miami, Florida

Professor Nezih Pala, Major Professor

Semiconductor nanostructures have attracted considerable research interest due to their unique physical and chemical properties at nanoscale which open new frontiers for applications in electronics and sensing. Zinc oxide nanostructures with a wide range of applications, especially in optoelectronic devices and bio sensing, have been the focus of research over the past few decades. However ZnO nanostructures have failed to penetrate the market as they were expected to, a few years ago. The two main reasons widely recognized as bottleneck for ZnO nanostructures are (1) Synthesis technique which is fast, economical, and environmentally benign which would allow the growth on arbitrary substrates and (2) Difficulty in producing stable p-type doping. The main objective of this research work is to address these two bottlenecks and find a solution that is inexpensive, environmentally benign and CMOS compatible. To achieve this, we developed a Sonochemical method to synthesize 1D ZnO Nanorods, core-shell nanorods, 2D nanowalls and nanoflakes on arbitrary substrates which is a rapid, inexpensive, CMOS compatible and environmentally benign method and allows us to grow ZnO nanostructures on any

arbitrary substrate at ambient conditions while most other popular methods used are either very slow or involve extreme conditions such as high temperatures and low pressure.

A stable, reproducible p-type doping in ZnO is one of the most sought out application in the field of optoelectronics. Here in this project, we doped ZnO nanostructures using sonochemical method to achieve a stable and reproducible doping in ZnO. We have fabricated a homogeneous ZnO radial p-n junction by growing a p-type shell around an n-type core in a controlled way using the sonochemical synthesis method to realize ZnO homogeneous core-shell radial p-n junction for UV detection.

ZnO has a wide range of applications from sensing to energy harvesting. In this work, we demonstrate the successful fabrication of an electrochemical immunosensor using ZnO nanoflakes to detect Cortisol and compare their performance with that of ZnO nanorods. We have explored the use of ZnO nanorods in energy harvesting in the form of Dye Sensitized Solar Cells (DSSC) and Perovskite Solar Cells.

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LIST OF ACRONYMS & ABBREVATIONS

MEMS	Microelectromechanical Systems
IEP	Isoelectric point
LED	Light emitting diode
CVD	Chemical Vapor deposition
1D	one dimensional
2D	Two dimensional
NWs	Nanowalls
NFs	Nanoflakes
NRs	Nanorods
CMOS	Complementary metal-oxide semiconductor
ITO	Indium Tin oxide
FTO	Fluorine doped Tin oxide
DSSCs	Dye sensitized solar cells
PSCs	Perovskite solar cells
FESEM	Field emission scanning electron microscope
TEM	Transmission electron microscope
AFM	Atomic force microscope
HRTEM	High resolution transmission electron microscope
PL	Photo luminescence
XRD	X- ray diffraction
CV	Cyclic Voltammetry
EIS	electrochemical impedance spectroscopy
ELISA	enzyme-linked immunosorbent assay

UV	Ultra violet		
PVD	Physical vapor deposition		
MOCVD	Metal organic chemical vapor deposition		
MBE	Molecular beam epitaxy		
PLD	Pulsed laser deposition		
SCCM	Standard Cubic Centimeters per Minute		
ZNH	Zinc Nitrate Hexahydrate		
ZAD	Zinc Acetate dehydrate		
TMD	Transition metal dichalcogenides		
EDS	Energy dispersive X- ray spectroscopy		
SAED	Selective area electron diffraction		
НМТ	Hexamethylene tetramine		
DI	Deionized		
PET	Polyethylene terephthalate		
LO	longitudinal		
IPCE	Incident photon to charge carrier efficiency		
FIB	Focused ion beam		
SCLC	Space- charge limited current		
AlN	Aluminum Nitrate		
Al:ZnO	Aluminum doped ZnO		
ТСО	Transparent conductive oxide		
LD	Laser diode		
PD	Photodetector		
P:ZnO nanoshell	Phosphorous doped ZnO nanoshell		
ADP	Ammonium dihydrogen phosphate		
n-ZnO	n type ZnO		

PMMA	Poly (methyl methacrylate)		
RIE	Reactive ion etching		
BSA	Bovine saline albumin		
NEP	Noise equivalent power		
MGJ	MoS ₂ /graphene junction		
FET	Field effect transistor		
SBJ	Schottky barrier junction		
SBH	Schottky barrier height		
MGPD	MoS ₂ /graphene photodetector		

CHAPTER 1

Introduction

1.1 Motivation and Goals

Semiconductor nanostructures have attracted considerable research interest due to their unique physical and chemical properties at the nanoscale, which opens new frontiers for applications in electronics and sensing. ZnO is one of the most important semiconducting material in (II-VI) group. ZnO possess a wide bandgap of 3.37 eV and a high exciton energy of 60 meV, which attracted the attention of scientific community dating back to 1935 especially in the short wavelength optoelectronics. [1] The lack of center of symmetry and strong electromechanical coupling makes ZnO piezoelectric and pyroelectric which was made use in MEMS as transducers. Apart from the above applications, ZnO is also an important material in bioelectronics owing to its high biocompatibility, less toxicity and high IEP value. [2, 3]

Zinc Oxide nanostructures with wide range of applications especially in optoelectronic devices and bio sensing have been the focus of research over the past few decades. There were numerous reports that illustrated the use of ZnO nanostructures in various applications such as pressure sensors, [4] biosensors, [5] gas sensors, [6] light emitting diodes (LEDs), [7] displays, [8] transparent electrodes, [9] and solar cells [10] to name a few. Another important feature of ZnO that has fuelled the interest in ZnO is the availability of diverse family of ZnO nanostructures [11, 12, 13] such as nanorods, [14] nanobelts, [15] nanoflowers, [16] nanocombs, [17] and nanotetrapods [18] to name a few.

Despite the immense potential and the many advantages they bring, devices based on ZnO nanostructures have failed to attain the success similar to other materials such as Titanium dioxide (TiO₂) or Gallium nitride (GaN) as was expected to, a few years ago. The two main reasons widely recognized as bottleneck for the development of devices based on ZnO nanostructures are (1) Synthesis technique which is fast, economical and environmentally benign which would allow the growth on arbitrary substrates (2) Difficulty in producing stable p-type doping. The main objective of this research work is to address these bottlenecks and achieve improved performance in ZnO devices.

Research spanning several decades has yielded several synthesis techniques to grow ZnO nanostructures with precise control over morphology and aspect ratios. Among all, the most widely used techniques are based on vapor transport process in which Zinc (Zn) and Oxygen vapor are transported and then allowed to react at the desired site for the formation of Zinc Oxide (ZnO) nanostructures. Formation of 'Zn' in vapor form requires high temperature, precise pressure control and ratio of 'Zn' to oxygen vapor which has a significant effect on the morphology of the nanostructures. Chemical vapor deposition (CVD) technique, which is a widely used vapor transport method uses ZnO and graphite powders to generate 'Zn' and oxygen vapors at temperatures from 800 °C to 1000 °C to grow highly crystalline nanostructures. However, the need for extreme growth conditions makes it an impractical solution for growth on flexible, plastic substrates. Chemical synthesis methods such as solvothermal methods, sol-gel method, electrodeposition technique and microwave assisted growth to name a few have been hugely popular methods as they can be performed at ambient conditions without use of many specialized equipment,

at the same time having control over the size and morphology of the nanostructures by varying the parameters such as pH, growth time and temperature. However while the CVD technique is limited by the extreme growth conditions which increase the cost of synthesis and places a serious limitation on the choice of substrate, chemical methods are largely limited by their long synthesis times and high level of hydrogen (H) and Oxygen (O) impurities in the as grown nanostructures. This calls for development of a comprehensive synthesis technique to synthesize ZnO nanostructures at ambient conditions on arbitrary substrates which can facilitate the use of flexible substrates for wearable electronics and minimize the defects. In this dissertation we developed a Sonochemical method to synthesize 1D ZnO Nanorods (NRs), ZnO homojunction and heterojunction core-shell nanostructures, 2D Nanowalls (NWs) and Nanoflakes (NFs) on arbitrary substrates. Sonochemical method is rapid, inexpensive, CMOS compatible and environmentally benign method which allows us to grow ZnO nanostructures on any arbitrary substrate at ambient conditions while most other popular methods used are either very slow or involve extreme conditions such as high temperatures and low pressure. ZnO with its direct band gap of 3.37 eV and high exciton energy (60 meV) is an important material in optoelectronic applications, however a low intrinsic electron concentration which results in low conductivity meant that its application in microelectronics industry is limited. Assembling or direct growth of ZnO on other materials with high conductivity is one way of improving the properties of the material yet retaining most of the optical and material properties of ZnO. Using Sonochemical method we were able to grow ZnO nanorods on various substrates including Gold, Aluminum, Iron and Graphene. Another popular method to increase the carrier concentration in semiconductors is doping. Group (iii) metals such as

Boron (B), Aluminum (Al), Gallium (Ga) and Indium (In) can be used to successfully dope ZnO to increase the electron concentration. We present for the first time n-type doping of ZnO nanorods with 'Al' using Sonochemical technique.

Despite its many advantages over other direct bandgap semiconductors, the applications of ZnO in optoelectronics is hampered by the lack of stable p-type doping due to self-compensation by donor like native defects. This is especially prominent in a low temperature, solution method growth process. A stable, reproducible p-type doping in ZnO is one of the most sought out application in the field of optoelectronics. The sonochemical synthesis method offers a major advantage over other solution methods as the amount of self-compensating defects introduced in the sonochemical process is very small compared to other solution methods which should significantly improve the chances of achieving a stable p-type doping in ZnO which has for so long eluded researchers. Here, in this dissertation we focus on fabricating a homogeneous ZnO radial p-n junction by growing a p-type shell around an n-type core in a controlled way using sonochemical synthesis. Coreshell hybrid nanostructures have inspired various applications as they help improve the performance of the native nanostructure. Core-shell structures were used to increase the functionality, improve the thermal and chemical stability, increase storage capacity and improve the efficiency and responsivity of optoelectronic devices by minimizing the recombination losses and improved optical properties. In this research we used the sonochemical technique to synthesize ZnO-ZnO homojunction core-shell nanorods and ZnO-ZnS heterojunction Core-shell nanorods.

We have explored the as synthesized nanostructures in various applications such biosensors for point of care Cortisol detection, UV photodetectors and energy harvesting devices such as Dye Sensitized Solar Cells (DSSCs) and Perovskite Solar Cells (PSCs). Apart from ZnO we have also explored other 2 D materials such as graphene and MoS₂ which have attracted the interest of research and industry lately owing to their extraordinary electronic properties such as ballistic transport over 0.4 μ m and extremely high carrier mobility at room temperature in graphene, and high luminescence from monolayer MoS₂. In this study we have fabricated and characterized a MoS₂/ graphene Schottky junction photodetector with wide spectral response.

1.2 Scope of this Dissertation

This dissertation describes sonochemical synthesis, characterization and applications of ZnO nanostructures. Chapter 2 to 5 discuss the synthesis and characterization of ZnO nanostructures in various morphologies while chapter 5 to chapter 7 describes in detail the applications and advantages of ZnO nanostructures. Chapter 8 provides a summary of the dissertation and proposed future work.

Chapter 2, gives the background of ZnO and describes its properties in detail. ZnO crystal structure and resulting mechanical, electrical and optical properties were listed. Various ZnO morphologies and synthesis techniques with their merits and de-merits were explained. This chapter also describes the origin of n–type doping in ZnO and the hardship in achieving p–type doping. The various attempts at achieving p-type doping in ZnO nanostructures were reviewed.

In chapter 3, the theory of sonochemistry and the method of applying sonochemistry to synthesize ZnO nanorods, ZnO-ZnS core-shell nanorods, ZnO Nanowalls and ZnO Nanoflakes at INSYST lab were discussed in detail. The morphological and crystallographic characterization of the nanostructures using, Field emission scanning electron microscope (FESEM), Transmission electron microscope (TEM), Atomic force microscopy (AFM), High resolution transmission electron microscopy (HRTEM), Photo luminescence (PL), X- ray diffraction (XRD) and Raman spectrum were presented. Electrical characterization using, four point probe station, Hall measurements were presented. The optical characteristics of the nanostructures studied by measuring the transmittance, quantum efficiency and photo gain was presented in detail.

Chapter 4 discusses the theory of doping in ZnO. The role of defects in natural ntype doping is presented. The defects also play a major role in limiting the p- type doping in ZnO. This chapters presents the various attempts at achieving p-type doping in ZnO using IA group and VA group elements. In this chapter, we also present the sonochemical approach employed for the first time to dope ZnO with Aluminum (Al) and Phosphorous (P) to achieve n-type doping and p-type doping respectively. The fabrication of a homogeneous p- n junction and its characterization was presented in detail.

Chapter 5 discusses the role of ZnO in biosensing. We have used ZnO NRs and ZnO NFs as immobilizing agents to detect cortisol. The as synthesized anodes for immunosensing were characterized using cyclic voltammetry (CV) and electrochemical impedance spectroscopy (EIS). The results were compared with standard ELISA results and found that the as synthesized cortisol sensor is 100 times sensitive than standard methods. It was also observed that the electron transfer rates on ZnO NFs is much higher

than ZnO NRs. Further in chapter 6, we present the successful attempt to fabricate Dye Sensitized Solar Cells (DSSC) with high open circuit voltage and discuss the difficulties in attaining high short circuit current in perovskite solar cells. We propose and provide initial results for fabrication of ZnO nanorods by using ZnO nanoflakes as seed layer in order to reduce the recombination losses. In chapter 7 we present the use of doped ZnO nanorods in P:ZnO nanoshell/ n-ZnO NR radial p-n junction configuration to detect UV radiation. We also explore next generation materials such as 2 D MoS₂ and graphene. We present the first p- MoS₂/ graphene Schottky junction photo detector, which exhibits a wide spectrum due to the use of both band to band edge detection and internal photo emission over the Schottky junction. In chapter 8, we summarize the contribution of this dissertation and propose two new devices, (1) Flexible and transparent LEDs and (2) Flexible perovskite solar cells.

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CHAPTER 2

Background and State of the Art

2.1 **Properties of Zinc Oxide**

Zinc Oxide (ZnO) is one of the most important common and widely studied semiconducting material in (II-VI) group. ZnO possess a wide bandgap of 3.37 eV and a high exciton energy of 60 meV which attracted the attention of scientific community dating back to 1935 especially in the short wavelength optoelectronics. [1] Zinc oxide thin films have traditionally been used as transparent electrodes and UV photodetectors. [2] The lack of center of symmetry and strong electromechanical coupling makes ZnO piezoelectric and pyroelectric which was made use in MEMS as transducers. Apart from the above applications, ZnO is also an important material in bioelectronics owing to its high biocompatibility, less toxicity and high isoelectric point (IEP) value. [3, 4]

The last few decades saw an unprecedented advancement in nanotechnology where, the electronic and mechanical properties depend greatly on morphology and the dimensions of the material. This made ZnO one of the most studied material owing to the ease with which ZnO nanostructures can be synthesized in diverse morphologies and controlled dimensions. In the past, researchers were successful in developing techniques to synthesize ZnO nanostructures in 1D nanorods, nanowires, nanotubes, 2D nanobelts, nanocombs and many other morphologies which showed remarkable features in fields of sensing and field emission. [5]

2.1.1 Crystal Structure, Lattice Parameters & Mechanical Properties



Figure 2.1 (a, b) Hexagonal wurtzite structure of ZnO crystal [6]

Even though ZnO like other II-VI semiconductors occurs in three crystal structures namely zinc blend, rock salt and wurtzite, naturally occurring ZnO is found to be thermodynamically stable in wurtzite hexagonal structure with a space group of P6₃mc. The lattice constants of ZnO measured by X-ray diffraction are a= 3.2426 (+/-) 0.0001 Å, b= 3.2426 (+/-) 0.0003 Å, 3.2426 (+/-) 0.0001 Å. [7] Wurtzite ZnO structure can be described as alternating planes composed of a zinc ion surrounded tetrahedrally by four oxygen ions and vice versa. These planes are stacked upon each other in c-direction. This tetrahedral coordination of Zn²⁺ ions and O²⁻ ions results in non-central symmetric structure which makes ZnO piezoelectric and pyroelectric. [8] Also the oppositely charged ions along the basal plane results in normal dipole moment and spontaneous polarization along

the c-axis results in a lack of symmetry in the wurtzite structure which enhances the piezoelectric nature of ZnO. [9] The typical values of piezoelectric and pyroelectric coefficients of ZnO thin film are 11.0 pC/N and -9.4 x 10^{-6} C/m²K respectively. This property has been widely used in making ZnO cantilevers and pressure sensors. [10, 11]

2.1.2 Electrical and Optical Properties

Zinc Oxide is a direct bandgap semiconductor material with a bandgap of 3.37 eV. This large bandgap which is near ultraviolent light (UV) range means ZnO devices have higher breakdown voltage, can sustain large electric fields and produce lower electrical noise. The large band gap also means ZnO is transparent to visible light. This property of ZnO has made it one of the most researched material for transparent electrodes as a replacement for the widely used indium tin oxide (ITO) and titanium dioxide (TiO₂). ZnO possesses a high exciton binding energy of 60 meV which is 2.4 times that of thermal energy at room temperature. This paves way for intense near-band-edge excitonic emission at room temperatures which was suggested to be the mechanism behind ZnO lasing. [12] Another important feature of ZnO is its surface sensitivity. The conductivity of ZnO is very sensitive to surface adsorbents, this coupled with its high iso-electric point (IEP) of 9.5 [4] and biocompatibility has made ZnO a highly desirable and exploited material in electrochemical biosensing and gas sensing. Other ZnO parameters and properties such as dielectric constant, carrier concentration etc. are listed in table 1.

Static dielectric constant	8.656		
Refractive Index	2.008, 2.029		
Energy Gap	3.4 eV		
Intrinsic carrier concentration	$<10^{6} \text{ cm}^{-3}$ (max n-type doping $> 10^{20} \text{ cm}^{-3}$		
	Electrons; max p-type doping > 10 ⁻⁷ cm ⁻⁵		
	holes)		
Exciton binding energy	60 meV		
Electron effective mass	0.24		
Electron Hall mobility at 300 K for	$200 \text{ cm}^2/\text{Vs}$		
low n-type conductivity			
Hole effective mass	0.59		
Hole Hall mobility at 300 K for low	$5-50 \text{ cm}^2/\text{Vs}$		
p-type conductivity			

Table 2.1. Electrical and optical properties of bulk ZnO [13]

2.2 Defects and Doping in ZnO

ZnO is naturally n-type, the reason being still debated. However the most widely accepted reason is the presence of unintentional intrinsic defects such as oxygen (O_v) and Zn (Zn_v) vacancies or Oxygen (O_i) and Zinc interstitials (Zn_i) or antisites such as Oxygen (Zn_o) and Zinc (O_{zn}) antisites [14]. Beside Oxygen and Zinc defects, doping of hydrogen donor which can be easily incorporated during the growth process can be a major reason for the unintentional n-type doing observed in ZnO. The average electron concentration in ZnO films is approximately 10^{21} cm⁻³. [13] Figure 2.2 shows energy levels of various defects present in the ZnO band gap.

E _c	V ²⁻ 56 eV	V _o — 0.05eV Zn _i + 0.5eV	Zn _i	(H _i + 0.03-0.05eV
,	V _{Zn}	V _o 1.62 eV, 2.3 eV	0 _i — 2.04 eV	0 _{zn}	V _o Zn _i — 2.16 eV
2.5	55 eV	V _o ⁺ 2.0 eV, 3.1 eV	0, — 2.96 eV		

Figure 2.2. Energy levels of various defects in ZnO [15]

This unintentional naturally occurring n-type conductivity in ZnO is advantageous in various applications such as transparent electrodes, but also forms a bottle neck in achieving p-type doping in ZnO as they self-compensate the p-type dopants. Several groups have reported partial success in doping ZnO with group-I and group-V elements as dopants. [16, 17, 18] However a successful and reliable method is yet to be realized. Here in this research work, we have successfully doped ZnO nonoshells with phosphorous using sonochemical synthesis.

2.3 ZnO Nanostructures and Synthesis Techniques

During the past 15 years, nanotechnology has progressed at an unprecedented phase. By virtue of their unique properties that differ significantly from their bulk counterparts, nanomaterials offer novel applications to the field of science and technology. As the size

of a material reduces to nanoscale, the movement and degree of freedom of electrons (and photons) is confined to a limited space. This quantum confinement alters the physical and chemical properties of the material compared to its bulk counterpart. At this scale, the morphology of the material plays a vital role in defining the properties of the material. The asymmetrical structure of the ZnO can induce growth in different growth directions depending on the growth conditions. This has encouraged researchers around the world to develop diverse morphologies of ZnO nanostructures by controlling the growth kinetics. Nanostructures such as nanorods, [19] nanowires, [20] nanowhiskers, [21] nanowalls, [22] nanobelts, [23] nanotetropods, [24] nanocombs [25] are few of the many morphologies of ZnO nanostructures that have been reported over the past decade using various synthesis techniques such as chemical vapor deposition (CVD), [26] physical vapor deposition (PVD), [20] metal organic chemical vapor deposition (MOCVD), [27] molecular beam epitaxy (MBE), [28] pulsed laser deposition (PLD) [29] and wet chemical process such as hydrothermal method, [30] electrodeposition method, [31] microwave irradiation assisted growth [32] and sol-gel method [33] to name a few. Among these techniques, CVD method and hydrothermal methods which attracted the interest of researchers have been widely reported over years. These techniques were studied in detail come with their own drawbacks.

CVD technique is performed by allowing ZnO powder to react with carbon (C) powder at high temperatures (800 °C to 900 °C) in a furnace. The transformation of ZnO powder and carbon at the elevated temperature in to Zinc (Zn) and Carbon dioxide (CO₂) gas forms the basis of CVD technique. These gases are carried by a carrier gas (Argon) to
a substrate coated with a catalyst (Au), and placed at a lower temperature in the same furnace. At lower temperature 'Zn' gas oxidizes back in to ZnO. The nucleation of the ZnO nanocrystal happens at the tip of the catalyst and develops in to nanostructures. The morphology of the nanostructures depends on various factors such as gas flow rate, temperature and pressure. Figure 2.3 (a) shows the mechanism of growth via CVD technique and figure 2.3 (b) shows the CVD set up at INSYST lab. Figure 2.3 (c) shows the ZnO nanorods grown on Si substrate using CVD technique where substrate is placed at 500 °C. In this process a catalyst (commonly Gold (Au)) is deposited on Si substrate. This catalyst at high temperature forms an alloy with the source material. When the concentration of the semiconductor reaches the upper solubility limit, nanowire nucleation commences. ZnO nanowires were synthesized in low pressure chemical vapor deposition system. In this method, equal amounts of Carbon (C) and ZnO powder are well mixed and used as the source material. The substrate is gold sputtered followed by annealing at 400 ^oC. The whole process is performed in a 3 zone furnace as shown in figure 2.3 (a). The source is maintained at 800 $^{\circ}$ C – 900 $^{\circ}$ C while the substrate is placed around 500 $^{\circ}$ C. The whole furnace is maintained at 63 Torr pressure. A constant Argon gas flow of 100 SCCM is maintained. The process takes 60 mins.



Figure 2.3 (a) Growth process in VLS method (b) In house 3 zone CVD furnace (c) ZnO nanorods grown via CVD method on Si substrate.

Various morphologies such as rods, whiskers, tripods and 2D nanobelts synthesized using this process have been reported. CVD process gives a very high quality single crystalline ZnO nanostructures but requires high temperatures (800 °C- 900 °C) which limits the choice of substrates that can be used. Also the need for catalyst can induce impurities in to ZnO nanostructures. Therefore this method cannot be used to grow ZnO nanostructures on flexible organic substrates for future flexible and wearable electronics.

Hydrothermal process is a chemical process, in which Zinc ions (Zn^{2+}) derived from precursors such as Zinc Nitrate Hexahydrate (ZNH) and Zinc Acetate Dehydrate react with OH⁻ ions from water (H₂O) and other alkaline solutions to form ZnO nanostructures. Hydrothermal method requires low temperatures from 100 °C to 180 °C and no additional catalysts making it a suitable method for organic and flexible substrates and requires simple setup making it an inexpensive process. However this method takes longer times up to 12-24 hours for the process to complete. The low temperature involved in the process have been known to result in higher defects, especially hydrogen defects when compared with high temperature CVD process. To be able to be a part of practical low cost electronic applications, these techniques are either not cost effective or require conditions too harsh to give any flexibility for the growth on desired substrates. In this dissertation, sonochemical methods which involves the use of high intensity ultrasonic irradiation of an aqueous solution was developed as an alternative method for the synthesis of ZnO nanowires.

Recently, there has been a lot of interest in the synthesis and study of 2D inorganic structures after the discovery of graphene. 2D Graphene has revolutionized nanoelectronics like no other material in the new century owing to its remarkable mechanical, electrical and optical properties. Graphene exhibits remarkably high mobility with a maximum reported mobility of $\sim 200, 000 \text{ cm}^2 \text{V}^{-1}\text{s}^{-1}$. [34] This has the potential to take the field of electronics and communications in to unchartered territories. However the lack of band gap has been the major obstacle in achieving the desired goals. Integrating other 2D semiconductor materials with graphene can bring best of both the worlds. Researchers around the world have dedicated their efforts in realizing 2D counterparts of many bulk semiconducting materials and integrate them with graphene. Recently, 2D atomic layers of semiconductor transition metal dichalcogenides (sTMDs) such as MoS₂, WSe₂ were successfully realized and explored for various optoelectronic applications. [35,36] These structures have also been assembled with graphene to form 2D hybrid materials which improve the functioning of both graphene and other 2D materials. [3,38] In this study, we present the fabrication and junction properties of MoS₂/ Graphene Schottky junction photodetector.

2.4 References

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CHAPTER 3

SONOCHEMICAL SYNTHESIS OF ZINC OXIDE NANOSTRUCTURES

3.1 Theory of Sonochemistry

Sonochemistry is a study of chemical reactions and crystallization induced in a solution after application of ultra sound. Introduction of ultrasound in to solution creates regions inside the solution with high and low pressure regions due to periodic compression and expansion. [1] This change in pressure initiates a process called acoustic cavitation, in which air bubbles are formed in the solution and proceeds to grow in size before they collapse releasing a large amount of energy in to its immediate surroundings. Air molecules dissolved in the solution diffuse to form bubbles at the low pressure cycle. During the next cycle, the bubble is subjected to immense amount of external pressure which compresses the bubble. This process of bubble growth and compression continues until it reaches a point where the external pressure is larger than the surface energy of the bubble at which point the bubble collapses releasing a large amount of energy and pressure in to its immediate vicinity. It was measured that temperatures as high as 5000 K and a pressure of 1000 atm can be reached during this process. The core of the bubble known as hot spot, features high energy collisions which generates an energy as high as 13 eV. [1,2] This property of acoustic cavitation and its effects has been the center of research for many years starting from Lord Rayleigh in year 1983 when he was assigned to determine the cause of the damages done to the propeller of HMS Daring, which shed light on both the constructive and destructive behavior of acoustic cavitation. He deduced that acoustic cavitation was the reason behind the damage caused to the propeller. He developed a

mathematical model for the formation of cavities in liquid and predicted the existence of hotspot inside the bubble. The sequential process of acoustic cavitation is shown in figure 3.1



Figure 3.1. Dynamics of cavitation bubble [2]

Rayleigh equation shown below gives relation between the size of the bubble, density of solution and pressure difference between compression and rarefaction cycle. The size and growth of the bubble reciprocates the change in pressure. [3]

$$R\ddot{R} + \frac{3}{2}(\dot{R})^2 = \frac{(p-p_0)}{\rho}$$
(1)

Rayleigh however, in his model did not take the viscosity and surface tension of the solvent in to account. Plesset in 1949 proposed a modified Rayleigh's equation to include the effect of viscosity and surface tension given by Raleigh-Plesset equation (2). [4]

$$R\ddot{R} + \frac{3}{2}(\dot{R})^2 + \frac{4\sigma\dot{R}}{R} + \frac{2S}{\rho R} = \frac{(p-p_0)}{\rho}$$
(2)

Where 'R' is the radius of the bubble and a function of time. ρ is density, σ is viscosity and *S* is the surface tension of the liquid. (*p*-*p*₀) is the pressure difference between compression and rarefaction cycle. This pressure difference can be calculated using simple wave equation given in equation (3).

$$P(x,t) = A\cos(\omega t - \beta x + \varphi)$$
(3)

where, 'A' is the amplitude, ω is 2π times the frequency of the sound wave, β is given by 2π divided by the wavelength and φ is the reference phase of the sound wave respectively. The amplitude and frequency determines the magnitude and rate of oscillation of the pressure field which is critical to the formation and sustaining of acoustic bubbles. The pressure decreases at rarefaction, such that the local temperature becomes sufficient for dissolved air bubbles and edge cavities to diffuse, a process known as rectified diffusion. During rectified diffusion, air cavities nucleate around a particle or conglomerate to form bigger cavity – a sub process known as nucleation. The diffusion process completes within a rarefaction cycle and leads to the oscillating growth of the bubble. After rarefaction, the new born bubble is squeezed and shrunk by the external pressure that tries to implode it. The bubble overcomes this high pressure cycle by adapting a suitable size. Depending on the applied amplitude and frequency, the size increases from few micrometers to hundreds of micrometer between the inception and final phase within few milliseconds. As the bubble reaches the compression at its peak, the external pressure dominates and the bubble implodes. The dying bubble is reduced instantaneously to ~ 10 th of its size before implosion; while the matter inside turns to plasma due to the violent compression. Filled with high energy ions diffused by the decomposing chemical solution, the energy inside the bubble can rise to as high as 13 eV from inter-ion collisions. This hotspot energy is then conferred to the ambient liquid along with the ions and chemicals upon bubble implosion which is used by sonochemists to facilitate chemical reactions and nanostructure growth, which is not achieved in normal solution process methods. Another attractive aspect of using sonochemistry in nanomaterial synthesis is the faster kinetics of chemical reactions, which reduces the time of reaction considerably.

Several factors affect the growth of nanomaterials using Sonochemistry. As lord Rayleigh found out, sonochemistry can be devastating if not controlled in an effective manner. In order to control the cavitation and chemical reactions in a solution, it is important for the solution to be a homogeneous mixture. In case of heterogeneous solution, the physical effects of cavitation becomes dominant. Small particles driven by shock wave undergo drastic structural transformation upon collision while, bigger particles participate in shock wave induced collisions which give rise to micro jets at the boundaries. These physical effects not only cause a severe damage to the equipment and particles used, but can also drive the process out of control. Therefore it is important to have a very homogeneous solution for the process. Apart from the homogeneity of the solution, other factors such as vapor pressure of the solvent and concentration of the precursors play a vital role in shaping the chemical reaction and determining the morphology of the nanoparticles that result from the cavitation process. [5] It was found that the viscosity and the surface tension of the solvent are other two factor which play a significant role in determining the size and density of the bubble which in turn, determines the intensity of the cavitation. The concentration of the solute also plays a major role in determining the morphology of the nanoparticles. [6] Another important player in the sonochemical process is frequency. The effect of change in frequency reflects on the pressure cycle which is crucial to the cavitation process. As the frequency increases the size of the bubbles become smaller and results in a uniform distribution. However, the energy buildup inside the bubble does not follow the increase in frequency. [7] It is crucial to find the optimum frequency range for enhanced cavitation process for a particular solution in order to achieve the desired chemical reaction. Other factors that affect the final outcome are sonication time, power and the substrate used. [8]

3.2 Sonochemical Synthesis of ZnO Nanostructures

At INSYST lab we make use of a simple, commercially available high intensity ultrasound set up which boasts a 750 W ultrasonic processor from Sonics and Systems as shown in figure 3.2. This system allows us to vary the amplitude of sonication and time of each cycle to a desired value, thereby giving a great control over the process. Using sonication method, we have successfully synthesized 1D ZnO NRs, 2D ZnO NWs, 2D ZnO NFs, ZnO-ZnO core-shell nanorods, ZnO-ZnS core-shell nanorods on arbitrary substrates. Extensive material characterization using, XRD, EDS, SAED and HRTEM has revealed that the nanostructures grown are single crystalline in nature and are oriented in (0002) direction.



Figure 3.2 Sonochemical Setup from Sonics and Systems

3.2.1 ZnO Nanorods

We used a two-step process for synthesis of ZnO nanorods on different substrates. In the first step, a uniform layer of ZnO seed was deposited on an arbitrary substrate by sonochemical method using 0.5 mM Zinc Acetate Dehydrate (ZAD) in 100 ml of isopropanol (IPA) as precursors. ZAD has been extensively studied for its ability to form seeds suitable for subsequent nanowire/nanorod growth on both polar and nonpolar substrates capable of tolerating 200 °C [9]. We modified this approach and utilized the effects of sonication to produce solution phase seeding and growth of nanorods in a solution at ambient conditions. Initially a 0.5 mM concentration of ZAD in isopropyl alcohol solution was prepared at room temperature by stirring with a magnetic stir bar at

350 rpm for 5 minutes. The substrates were cleaned with methanol, acetone and isopropanol before immersing in to the solution. The solution with immersed substrate was then irradiated using commercially available high intensity ultrasound setup in 15 minute cycles at 75% of the maximum amplitude (\sim 30 W cm⁻²) and 20 kHz frequency. The solution was refreshed after each cycle, the substrate is rinsed with DI water, dried with nitrogen and the process was then repeated as necessary to obtain a desired distribution and uniformity. During the whole process, the global temperature of the aqueous solution never exceeds \sim 70 °C which makes it suitable flexible plastic substrates. Figure 3.3 shows the SEM picture of ZnO seed layer on Si grown over 4 cycles of sonication with each cycle running for 15 minutes. The thickness and uniformity of the seed layer increases with the number of cycles while the diameter of seed increases with duration of each cycle.



Figure 3.3. Cross- sectional view of ZnO seed layer on Si substrate.

Following the seeding process, ZnO nanorods were synthesized by sonicating a mixture of 0.020 M Zinc Nitrate Hexahydrate (ZNH) and 0.020 M Hexamethylene tetramine (HMT) solution in 100 ml of deionized water (DI) at 50% amplitude (~20 W cm⁻²). The samples were rinsed with DI water and then dried in air after each cycle. When the solution containing ZNH and HMT in DI water is sonicated, OH, O₂, HO₂ and H radicals are produced. The OH⁻ ions are released by both water and HMT. The released OH⁻ ions promote the precipitation of ZnO from Zn(OH)₂. The detailed chemical reaction is shown below. [10]

 $(CH_2).6N_4 + 6H_2O \rightarrow 4NH_3 + 6HCHO$ $NH_3 + H_2O \rightarrow NH_4 + OH^ Zn(NO_3).6H_2O \rightarrow Zn^{2+} + 2NO^{3-}$ $Zn^{2+} + 2OH^- \rightarrow Zn(OH)_2$ $Zn(OH)_2 \rightarrow ZnO + H_2O$

These reactions take place extremely slow under normal room temperature. A typical process takes 24 hours even at concentration of precursors above 10 m mol. However under sonication, these reactions take place extremely fast which gives a growth rate as high as 10 nm/min. Apart from providing OH⁻ ions, HMT also plays a vital role in directional growth of ZnO nanorods. HMT a long chain polymer and a nonpolar chelating attach to the nonpolar facets of Zinc Oxide there by cutting the supply of further Zn²⁺ ions as shown in figure 3.4. This leaves only the polar (0001) face for epitaxial growth. [11, 12] The topography and morphology of the nanorods was studied using Scanning electron microscopy (SEM) and Atomic force microscopy (AFM), shows that ZnO nanorods grow in vertical direction.



Figure 3.4. HMT attached to non- polar facets of ZnO resulting in a directional growth of ZnO nanorod



Figure 3.5 shows the SEM images of Sonochemically grown ZnO nanorods on (a) Silicon (b) SiO_2/Si (c) Ni/Fe alloy (d) Teflon (f) Copper (g) Iron (h) ITO (i) Sapphire (j) Kaplan tape (k) cotton fiber (l) PET

To understand the effect of number of cycles and the duration of each cycle, ZnO NRs were grown by sonicating the solution with constant concentration over varying times and cycles. It was observed from the SEM pictures that the length of nanorods increases with the duration of each cycle as the rod length increases from 500 nm after 15 minutes of growth to 750 nm after 30 minutes of growth.



Figure 3.6 (a) ZnO nanorods after 15 minutes of growth (b) ZnO nanorods after 30 minutes of growth

The length also increases with increase in cycles as it increases from $0.36 \,\mu\text{m}$ for 5 cycles to 2 μm for 26 cycles. The increase in diameter with number of cycles is more conservative however, as it was observed that the diameter plateaus around 50 nm. The results of this study are plotted in figure 3.7

In summary, we have achieved control over the growth of ZnO nanorods as the density of growth can be increased with increase in number of cycles, while the duration of each cycle determines the length of the ZnO nanorods. ZnO nanorods as long as $3 \mu m$

can be achieved while keeping the diameter constant at 50 μ m which gives high surface to volume ratio.



Figure 3.7. Change of dimensions with growth cycles in ZnO nanorods

3.2.1.1 Properties of ZnO Nanorods

Photo Luminescence (PL), X-ray diffraction (XRD) and Transmission electron microscopy were used to study the morphology and crystallographic properties of ZnO nanorods. Figure 3.8b shows the PL of ZnO nanorods grown on ITO. It can be noted that the spectrum is comparable to the catalyzed growth on 'Ag' and 'Zn' coated substrates which have shown to result in high quality ZnO nanostructures. The intensity of 380 nm peak observed is significantly higher than the defect peak at 550 nm .[13] This suggests that sonochemical synthesis results in high quality, low defect nanorods unlike other solution methods.



Figure 3.8. (a) ZnO nanorods on ITO (b) PL of ZnO nanorods showing a significant peak at 380 nm.

Further structural analysis was performed by transmission electron microscopy (TEM). Figure 3.9 (a) shows the selected area diffraction pattern (SAED) which confirms the single crystalline wurtzite structure of the ZnO nanorods with a dominant growth direction along (0002) plane. The same can be confirmed from the HRTEM image shown in figure 3.9 (b). The HRTEM images show that the synthesized nanorods are structurally uniform and contain no defects such as dislocations and stacking with a measured lattice spacing of \sim 0.5 nm which corresponds to the d-spacing of (0002) crystal planes.

Figure 3.10 (a) shows the X-Ray diffraction pattern, which is a typical XRD pattern for hexagonal wurtzite ZnO nanorods. [14] Figure 3.10 (b) shows the Raman spectra

obtained from ZnO nanorods grown over a Si substrate using a 532 nm laser. We observed a dominant E_2 (high) phonon peak at 436 cm⁻¹ which is blue shifted by 3 cm⁻¹ when compared to the bulk ZnO. [15]



Figure 3.9 (a) SAED pattern of ZnO nanorod (b) HRTEM image of ZnO nanorod, both confirming the single crystal nature of the nanorod

The 'LO' phonon peaks observed in bulk material at 574 cm⁻¹ which corresponds to longitudinal phonon vibrations were not observed in 1D ZnO nanorods, as the incident light is perpendicular to the c-axis of wurtzite ZnO.



Figure 3.10 (a) XRD on ZnO nanorods showing a relatively high intensity (002) peak (b) Raman spectra of ZnO nanorods with E_2 peak at 436 cm⁻¹.

The optical transmission spectra was measured using IPCE measurement system from PV measurements and is shown in figure 3.11 (a). To measure the transmission spectra and the dependence of transmission on the density and thickness of ZnO nanorods grown, we prepared two samples with ZnO nanorods grown in one and two cycles respectively over a glass substrate. A sharp valley denoting absorption is observed on the curve at 368 nm corresponding to the 3.37 eV which agrees with the bandgap in ZnO nanorods reported in literature. [16] Figure 3.11 (b) shows the electrical response of ZnO nanorods in dark and under UV illumination.



Figure 3.11 (a) Optical transmittance of ZnO nanorods, with the dip at 368 nm, corresponding to the bandgap of ZnO nanorods (b) Current- Voltage measurements in dark and under UV light, ZnO nanorods show a 3 order magnitude gain in presence of UV light.

We observed a three order magnitude gain in current in the presence of UV light. This photo response is due to the fast carrier thermal excitation and trapping at the surface of ZnO nanorods. During the growth of ZnO nanorods, there is usually a high density of defects at the surface which traps holes and increases photocurrent response. When UV

light is present, the photon energy is larger than the band gap of ZnO which generates electron hole pairs. While the high density surface trapping states are leaving the material, the photo-generated holes are trapped along the surface. The photo generated electrons are left unpaired and therefore contribute to the rapid increase in photocurrent. This property of ZnO has been used in many applications such as UV photodetectors, UV light emitting diodes and UV lasers.

3.2.2 ZnO Core-Shell Nanorods

Since the start of the century, researchers saw an increased interest in core-shell nanostructures such as core-shell microspheres, [17] nanotubes, [18] nanobelts, [19] nanorods [20] etc. Core-shell structures have shown to enhance the properties of semiconductor materials such as increased functionality, improved thermal and chemical stability, modulated band gap, improved efficiency in solar cells, increased storage capacity, and controlled drug delivery. [21-23] In this dissertation, we have used sonochemical method to develop ZnO-ZnO homojunction nanorods and ZnO-ZnS heterojunction nanorods.

3.2.2.1 ZnO-ZnO Homojunction Nanorod

Long has researchers controlled the growth direction of ZnO nanocrystals and growth kinetics of ZnO facets to synthesize diverse morphologies. Surfactants such as citrate ions were used to cap the basal plane of ZnO and thereby control the growth in 'c' direction. However this method introduces unnecessary impurities in to the ZnO crystal, moreover it

also introduces surface defects which increase the surface recombination when used in opto-electronic devices. In this dissertation, we have developed a sonochemical method to synthesize core-shell nanorods by varying the concentration of HMT (a surfactant in itself) to enhance the growth in planes other than c-direction. Core-shell nanorods were also synthesized by using citrate ions as surfactants for comparison purpose. Figure 3.12 (a-c) shows the SEM images of ZnO nanorods before and after the growth of ZnO shell using both citrate ions and HMT. Figure 3.13 (c) shows the TEM picture of ZnO- ZnO coresshell nanorod grown by varying HMT concentration.



Figure 3.12 (a) ZnO nanorods grown by sonochemical method with an average diameter of 55 nm (b) ZnO- ZnO core- shell nanorods grown by varying the ratio of ZNH and HMT with an average 60 nm shell thickness (c) Shell structure grown around ZnO rod by assistance of sodium citrate as surfactant



Figure 3.13 (a) ZnO nanorods grown by sonochemical method with an average diameter of 55 nm (b) ZnO- ZnO core- shell nanorods grown by varying the ratio of ZNH and HMT (c) TEM image of ZnO- ZnO core shell nanorod, clearly showing a shell encapsulating 60 nm thick nanorod

From both SEM and TEM pictures we observe that the diameter of the ZnO core is 50 nm-60 nm and the thickness of shell to be 30 nm. It can be observed that the latter case forms a much smoother and continuous interface between the core and shell. This structure will be used in formation of p-n homojunction discussed in the proceeding chapters.

3.2.2.2 ZnO-ZnS Heterojunction Nanorod

ZnS a wide band gap semiconductor with a band gap of 3.7 eV is a widely used luminescence material with applications in sensors and lasers. [24] Modifying the surface of ZnO by growing a ZnS shell around ZnO core is a well-recognized method to improve the functionality, chemical stability and optical response of the nanostructures, especially ZnO-ZnS core shell structures were used to improve the performance of solar cells where ZnS acts as a barrier between electrons and holes and thereby reducing the recombination loses. [25, 26] Many researchers have reported the synthesis of ZnO-ZnS core shell structures by sulfuration of ZnO core. However, there was no report known to have reported the growth of ZnO-ZnS core-shell nanorods using sonochemistry. We have for the first time synthesized ZnO-ZnS core shell nanorods using sonochemistry.

The ZnO core material was synthesized as described in the earlier section. To develop ZnS shell, a solution of 0.6 M sodium sulfide (Na₂S) solution was first prepared and then added drop wise to a solution of 0.02 M ZnH and 0.02 M HMT until a pH value of 6.0 - 6.15 is achieved. Already prepared substrate with ZnO nanorods is immersed in to

this solution and sonicated for 15 minutes. During this process the Sulfur replaces 'O' from the outer surface of ZnO forming ZnS shell. The chemical reactions taking place are:

 $Na_{2}^{S} + H_{2}^{O} \rightarrow NaHS + NaOH$ $NaHS + H_{2}^{O} \rightarrow H_{2}^{S} + NaOH$ $H_{2}^{S} + ZnO \rightarrow ZnS + H_{2}^{O}$ $NaHS + ZnO \rightarrow ZnS + NaOH$

Figure 3.14 (a-b) shows the SEM picture of ZnO- ZnS nanorods grown by sonicating for different times and figure 3.14 (c) shows the XRD data from ZnO-ZnS core shell nanorods. The peak at $2\Theta = 28^{\circ}$ indicates the formation of (111) plane of ZnS.



Figure 3.14 (a) ZnO- ZnS core-shell nanorods, shell grown for 10 minutes resulting in a thin shell (b) With longer process time the shell thickness increases depleting the ZnO core (c) XRD data on a ZnO-ZnS nanorods, a clear peak at 28° indicating the successful formation of ZnS

3.2.3 ZnO Nanowalls:

The same experimental procedure as above when performed over alumina (Al_2O_3) and aluminum (Al) coated substrates resulted in growth of ZnO Nanowalls instead. We later observed that there is no need for seed layer deposition as well. Figure 3.15 shows ZnO Nanowalls grown over 6 nm thick 'Al' over silicon substrate. The growth is largely substrate dependent as such growth was not observed on any other substrate.



Figure 3.15. ZnO Nanowalls grown over 6nm thick Aluminum coated Si substrate

To further understand the growth mechanism and kinetics of this process we conducted time dependent study. ZnO NWs were grown over Al coated Si substrate for 3 minutes, 6 minutes, 9 minutes, 12 minutes, 15 minutes, and 60 minutes. The growth kinetics and morphological changes were studied under FESEM. As shown here in figure 3.16 (a-f) ZnO NWs forms within first 9 minutes and by 15 minutes a complete ZnO NW network is formed. Further PL and Raman studies were used to study the growth mechanism. Figure 3.17 (a-b) shows the Photoluminescence (PL) and Raman spectroscopy of ZnO Nanowalls. The absorbance peak at 368 nm is observed as an indication of the inherent property of ZnO. The 378 nm PL peak is attributed to the free-exciton annihilation in ZnO. [27] Vaneusden et al. show that the large green emission is related to a deep-level or trap-state emission. They suggest that the green PL in ZnO is due to the recombination of electrons in singly occupied oxygen vacancies with photo-excited holes in the valence band. [28, 29]

Raman spectroscopy shows an emergence of a faint peak at 419 cm⁻¹ at 15 minutes which indicates the presence of $ZnAl_2O_4$.



Figure 3.16. Growth kinetics of formation of ZnO Nanowalls at (a) 3 minutes (b) 6 minutes (c) 9 minutes (d) 12 minutes (e) 15 minutes and (f) 60 minutes



Figure 3.17 (a) the PL peak at 378nm are a good indication of ZnO material. The visible peak is indicative of the defect levels commonly observed in ZnO. (b) Raman peak intensity for 3 minute intervals and at 60 minutes. We observe a faint peak at 419cm⁻¹ at 15 minutes and a more distinct peak at 660cm⁻¹ is observed for ZnAl₂O₄. The 520cm⁻¹ peak is attributed to Si, the 937cm⁻¹ is distinctive of Zn₂SiO₄. The incrementally increasing peak at 1053cm⁻¹ is attributed to dioxygen. The 438cm⁻¹ peak is attributed to ZnO

From the PL and Raman measurements a possible growth mechanism was proposed as shown below.

$$(CH_{2})6N_{4} + 6H_{2}O \rightarrow 4NH_{3} + 6HCHO$$

$$NH_{3+}H_{2}O \rightarrow NH + OH$$

$$Zn(NO_{3})6H_{2}O \rightarrow Zn^{2+} + 2NO^{2+}$$

$$Zn^{2+} + 2OH \rightarrow Zn(OH)_{2}$$

$$Zn(OH)_{2} \rightarrow ZnO + H_{2}O$$

$$Al + 2H_{2}O \rightarrow 1.5H_{2} + (AlO)OH$$

$$Zn(OH)_{2} + 2AlO(OH) \rightarrow ZnAl_{2}O_{4} + 2H_{2}O$$

The induced high pressure and temperature from the cavitation bubble imploding produces AlOOH during cavitation. The zinc/aluminum interface allows for zinc aluminate (ZnAl₂O₄) or "spinel" to form. [30] The reaction between Zn(OH)₂ and AlOOH allows for the formation of ZnAl₂O₄, this zinc aluminate has been synthesized in a hydrothermal solution at high temperatures. [31] HMTA and ZnAl₂O₄ would then support the vertical growth {0001} of ZnO Nanowalls. We find that the ZnO Nanowalls grow perpendicular to the alumina/spinel interface.

The crystal quality and orientation of the ZnO Nanowalls were studied by using selective area electron diffraction (SAED) using a 300 KV transmission electron microscope (TEM). Figure 3.18 (a-c) shows the lamella extracted using focused ion beam (FIB) tool along with TEM image and SAED pattern. To ensure a single layer of nanowalls was characterized, we extracted a thin lamella using focused ion beam (FIB) tool and measured SAED at the very edge of the sample. SAED patterns show that the nanowalls are single crystalline throughout the sample.



Figure 3.18 (a) SEM picture of the FIB prepared lamella of the ZnO Nanowalls on Al coated Si substrate (b) TEM image of Nanowalls (c) SAED pattern of ZnO NWs confirming the single crystalline nature of ZnO NWs

Current-voltage characteristics of ZnO nanowalls measured at dark and under UV

illumination are shown in Figure 3.19.





The measurements were taken on nanowalls synthesized on 10 nm thick patterned 'Al' electrodes on a Si substrate with a separation distance of \sim 50µm between the two tungsten probe tips. The grown ZnO nanowalls show a very low conductivity which was

also observed and reported by Lee et al. who grew ZnO nanowalls using low pressure MOVPE.[32] The reason for low conductivity even in the presence of ultraviolet (UV) light can be attributed to the formation of ZnAl₂O₄ during the growth of the nanowalls. It has been reported that ZnAl₂O₄ is lower in electrical conductivity than ZnO and other oxide mixtures. [33] The clear UV photoresponse is typical of ZnO material. Although a 0.1 W/cm² intensity UV light was exposed to the nanowall sample, we observe only one magnitude difference in the conductivity. A high photoconductive gain is usually reported for ZnO material. [34] We do not observe this in case of ZnO nanowalls.

3.2.4 ZnO Nanoflakes

We have synthesized extremely thin 2D ZnO nanoflakes on arbitrary substrates for the first time using sonochemical method. This method is a one step process which does not require a seed layer. This process was successfully tested on various substrates such as Si, PET and graphene. Growth of ZnO nanoflakes draws heavily upon the process used for nanorod growth. However no seed layer is required for this process and can be grown on most substrates available. The precursors used in this method are 0.2 M ZNH and 0.2 M HMT. A solution of 0.2 M ZNH and 0.2 M HMT was prepared at room temperature by stirring with a magnetic stirrer at 350 rpm for 5 minutes to ensure a homogeneous solution. The substrates cleaned with acetone, isopropanol and finally with DI water were immersed in the as prepared solution and irradiated using the same ultra sound set up used for ZnO nanorod growth. However the amplitude of sonication is set at 70% which results in ~ 30 W cm⁻² of energy. After the running a cycle for 15 minutes, the samples were cleaned with DI water and dried in air. For improving the density of growth, this process is repeated for

further cycles. This process results in extremely thin hexagonal shaped nano crystals which eventually self-assemble in to parallelogram shaped ZnO NFs. It was observed that the thickness of the nanoflakes increase with the duration of each cycle with the thinnest flake measuring as thin at 12 nm and the thickest around 200 nm.

The synthesis process of ZnO nanoflakes differs from nanorod growth only in terms of the concentration of the precursors primarily. The concentration of ZNH and HMT are increased 10-20 times and the sonication amplitude to 70%. The higher concentrations and the fast reaction rates achieved under intense cavitation causes HMT (was primarily used as shape inducing polymer in nanorod growth) to release a large amounts of OH^{-} and NH_{4}^{+} ions in a short period of time resulting in precipitation of Zn^{2+} ions at a phase much higher than the surface adhesion rate of HMT to the non-polar facets of ZnO. [35, 36] This results in growth of nonpolar planes of ZnO along with the (0001) plane, forming parallelogram shaped 2D ZnO nanoflakes. To further understand the crystallization and growth process, time dependent study of SEM images along with TEM and HRTEM was studied. Figure 3.20 (a-b) shows the SEM pictures of ZnO nanoflakes grown on Si substrate at 30 seconds and 60 seconds respectively. From the SEM images we can observe that the 2D ZnO nanoflakes grows initially in the form of thin sheets and then self assembles in to hexagonal crystal structures which further stabilizes in to parallelogram like structures. To further study the growth dynamics, samples were imaged at different time intervals of growth from 1 minute to 60 minutes. Figure 3.20 (b-e) shows the time evolution of the growth process for ZnO nanoflakes on Si substrate at 1, 3, 15 and 30 minutes respectively. Further SAED and HRTEM were used to identify the crystal edges of the ZnO nanoflakes as shown in figure 3.21. On the basis of SEM, HRTEM and SAED data we infered that the anisotropic growth of ZnO paves way for the observed morphology. This is explained in detail using figure 3.22. Initially, ZnO sheets are self- assembled in to hexagonal single crystalline structures with (0001) plane in c-direction. The crystal continues to grow in (01-10) direction which results in the disappearance of (01-1) plane and the resulting crystal takes the shape of more stable parallelogram. The longest edge is in (10-10) plane. It was observed that beyond a growth time, the crystal growth in lateral direction saturates beyond which, not much change in dimensions can be observed. Instead, new flakes start nucleating on the top of the old flakes which increases the thickness of the ZnO layer.



Figure 3.20 (a) ZnO nanoflakes grown over Si for 30 seconds (b) ZnO nanoflakes grown over Si for 60 seconds (c) ZnO nanoflakes grown over Si for 3 minutes (d) ZnO nanoflakes grown over Si for 30 minutes



Figure 3.21 (a) SAED pattern showing a single crystalline wurtzite structure oriented in <0001> direction and lateral growth direction perpendicular to (0110) (b) HRTEM at the edge of a flake with fringes almost parallel to the edge with a spacing of 2.74 Å, which is close to 2.82 Å of (1010) plane





Figure 3.23 summarizes the change of lateral dimensions and the thickness of the flakes by growth period. The thickness measurements were carried out using AFM, while lateral dimensions were measured using SEM. Sizes of several flakes have been averaged to have

statistically reliable data. The average thickness of the nanoflakes layer varies from 15 nm of single flake to 300 nm of a stack of several flakes for a growth time of 10 minutes and two cycles of 30 minutes each. However, the thickness of a single nanoflakes remains constant around 15-20 nm irrespective of the duration of each cycle. The lateral dimensions increases quickly form 0.76 μ m for growth time of 1 minute to 10 μ m at 15 minutes and remains the same as the time is increased. Therefore by controlling the time of each cycle and the number of cycles, the lateral dimensions, thickness and density of growth can be efficiently controlled using sonochemical method.



Figure 3.23 Evolution of length and thickness of the flakes by time showing that the lateral dimensions of the flakes increases rapidly to around 10 μ m at around 15 min and then stabilizes. The thickness of the single flakes remains constant around 20 nm

3.2.4.1 Properties of ZnO Nanoflakes

SAED, HRTEM, EDS, PL and Raman studies were conducted to study the crystallographic properties of the as grown ZnO nanoflakes. SAED pattern and HRTEM image confirm the

single- crystalline wurtzite structure of 2D ZnO nanoflakes oriented in [0001] direction. The HRTEM images are taken at the major edge of the flake show that the fringes are almost parallel to the edge with a lattice spacing 2.74 Å. This lattice spacing is a representative of (10-10) plane which corresponds to the longest edge. The SAED measurements indicate that the nanoflakes exhibit hexagonal wurtzite structure with a= 3.25 Å and c= 5.21 Å, oriented along (0001) plane. The photoluminescence (PL) spectra is shown in figure 3.24 (a) shows an emission line at 376 nm, which is characteristics of wurtzite ZnO. The peak intensity is relatively low due to extremely thin nature of the nanoflakes. The 376 nm PL peak is attributed to the free- exciton annihilation in ZnO. [37] The PL system used in this study upconverts the frequency from green light (532 nm) which causes the high- intensity stray green light, making it difficult to observe green PL emission. Raman spectra of ZnO nanoflakes is measured using 532 nm laser, is shown in figure 3.24 (b). Unlike in 1D ZnO nanorods, the prominent peak in 2D ZnO nanoflakes is not E_2 (high) but the highest intensity is observed at second order longitudinal phonon vibration (2LO) at 1050 cm⁻¹ which is due to the confinement in 'c' direction.



Figure 3.24 (a) Raman spectrum of ZnO on Si, a strong second order longitudinal phonon vibration peak characteristic of confinement in c direction is observed at 1050 cm⁻¹ (b) PL spectrum of the ZnO nanoflakes on graphene shows a peak at 376 nm.

In order to study the conductivity of the individual ZnO nanoflakes and the variation of conductance with the thickness, three samples Z-3, Z-5 and Z-15 were prepared by growing ZnO nanoflakes on Si/SiO₂ substrate with growth cycle of 3 minutes, 5 minutes and 15 minutes duration respectively. Electrical contacts were deposited using a focused ion beam (FIB) and I-V characteristics were measured using a probe station (The micromanipulator and results analyzed by semiconductor parameter analyzer, model no. Agilent 4156 A) attached with micromanipulator tip (model 7A-M) as shown in figure 3.25 (a). It can be observed that the I-V curves show slightly super linear behavior and the conductance increases with the increase in thickness of ZnO nanoflakes. Super linear behavior is attributed to space-charge limited current (SCLC), which is commonly observed in nanostructures of various materials including ZnO nanowires. [38] When the number of carriers injected from contacts is greater than the free carriers in the semiconductors, the current is limited by the buildup of injected space charge. Since SCLC is sensitive to the density of trap states and their energy depth, it can be prominently observed in 2D structures with large surface to volume ratio. Photoconductivity was measured on sample 3Z-15. 3Z-15 is prepared by running 3 cycles of 15 minutes each. This ensured a continuous film growth on Si/SiO_2 substrate. Figure 3.16 (b) shows the current -voltage curves for the said sample in dark and in presence of UV light. A 10 fold increase in conductance in the presence of UV light is observed, a characteristic of ZnO nanostructures caused by the fast carrier thermal excitation and trapping at the surface of ZnO nanoflakes. In the presence of UV light the photon energy is larger than the band gap of ZnO which generates electron hole pairs. The photo generated electrons are left unpaired

and therefore contribute to the rapid increase in photocurrent. This property of ZnO was observed in other nanostructures as well.



Figure 3.25 (a) Current–voltage (I-V) for samples of growth time 3, 5, 15 min show that the conductance of the flakes increases with the time of synthesis as the thickness increases and new layers are added with time (b) I-V for continuous fi lm of ZnO nanoflakes with growth time 15 min and three cycles show a 10-fold increase in conductivity in the presence of UV light

Graphene is a 2D carbon nanostructure, which has been studied extensively in recent years due to its remarkable mechanical, optical and electronic properties. Integrating graphene with other 2D materials so as to reap the benefits of both the materials has been the focus of many researches. Growing oxides over graphene can render additional functionality to graphene in electronic and optoelectronic applications as the hybrid material can have a band gap which is missing in graphene and at the same time exhibit high mobility. In this dissertation we used sonochemistry to grow ZnO NFs directly over graphene. For this purpose, a CVD grown monolayer graphene was transferred on to Si/SiO₂ substrate and was used as a substrate for ZnO nanoflakes growth. The resulting growth was compared with another sample prepared on Si/SiO₂ substrate under similar conditions. We observe
an increased density of growth on graphene as shown in figure 3.26. This increased growth can be attributed to the similar hexagonal atomic configuration in the "c" plane of both materials, which act as good template for the growth of hexagonal ZnO nanostructures. [39] Also the high conductivity of the graphene might have accelerated the growth of ZnO nanoflakes. Figure 3.27 shows the TEM, SAED and HRTEM pictures of ZnO nanoflakes on graphene. HRTEM and SAED confirms the single crystalline wurtzite structure oriented in (0001) plane. From HRTEM we observe that at the boundary of ZnO NF and graphene the nanoflake loses it's crystalline nature and is more amorphous. This is thermodynamically stable structure in nanomaterials. [40]



Figure 3.26 (a) ZnO nanoflakes on SiO₂ Substrate (b) graphene showing aggressive growth on graphene making the growth process highly selective

Ti/Au contacts were deposited on ZnO NF/ graphene sample to study the current-voltage characteristics in dark and under illumination. As shown in figure 3.28 (a) the conductance of ZnO NF/graphene is 1000 times higher than that of ZnO nanoflakes grown under similar conditions. ZnO nanoflakes/Graphene can be used as transparent electrodes provided they show a good transmittance along with the high conductance they showed. As shown in

figure 3.28 (b), the electrodes show a transmittance as high as 80% in the visible spectrum and a strong absorbance in UV spectrum.



Figure 3.27 (a) ZnO nanoflakes/ graphene on a TEM grid (b) TEM image of Individual flakes on graphene (c) SAED image of ZnO NF/graphene showing single crystal nature of ZnO NF (d) HRTEM image at the ZnO NF and graphene interface, ZnO nanoflake shows a single crystalline nature but at the interface where amorphous like behavior is observed.



Figure 3.28 (a) Electrical conductivity comparison in the presence of graphene and without graphene shows that the conductivity increases for the hybrid structure ZnO–graphene by more than 10 times when compared with just ZnO. (b) Comparison of transmission for graphene and ZnO–graphene hybrid structure

3.3 Conclusions

In conclusion, we have successfully developed sonochemical synthesis technique to synthesize, high quality single crystalline ZnO nanostructures such as 1D ZnO nanorods, ZnO-ZnO core-shell nanorods, ZnO-ZnS core-shell nanorods, 2D ZnO nanowalls and 2D ZnO nanoflakes on arbitrary substrates. The morphology and crystal structure of the nanostructures studied using SEM, TEM, XRD and Raman spectrum has shown that all the nanostructures were highly single crystalline in structure with very low levels of impurities. The electrical and optical properties measured indicate that the as grown nanostructures have applications in transparent electrodes and photodetectors.

3.4 References

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CHAPTER 4

Doping of Zinc Oxide

4.1 Al-doped ZnO

Despite the many promising features of ZnO discussed earlier, a low intrinsic electron concentration in undoped ZnO nanorods limits its use in various optoelectronic applications. Group 2A and 3A elements were used to dope ZnO in order to increase the carrier concentration. Numerous reports were published over the years to show an improved performance of ZnO owing to doping by Ga, [1] In, [2] Al, [3-5] Mn, [6] Mg. [7] Fortuna *et.al* have doped ZnO using Ga to achieve carrier concentration of 1.3×10^{21} cm⁻³, when ZnO nanobelts were doped with 'In' Jiansheng Jie et. al have observed that the carrier concentration increased to 7 x 10¹⁹ cm⁻³ .[8] Similarly Sasanka Deka and Vinod Kumar Etcheri have found an improved electrical response from ZnO nanparticles doped with Mn and Mg respectively. [9, 10] Similarly 'Al' was used to dope ZnO by various groups and found improved electrical properties without effecting its optical and crystalline properties to a greater degree which makes 'Al' a highly desirable dopant. [11] However, the high oxidation reactivity of 'Al' sources and the difference in vapor pressure between 'Al' and 'Zn' present a huge challenge for researchers to attain successful doping using conventional methods such as thermal evaporation and vapor-liquid-solid method. Various studies have shown to overcome these obstacles and achieve successful 'Al' doping using thermal evaporation [12] and vapor-liquid-solid method [13] but these methods require extensive setup, sophisticated equipment and high temperatures. In contrast, solution methods which use chemical approaches requires very simple equipment and can be performed at very low temperatures which make them more effective method to synthesize ZnO nanostructures on a large scale. In this dissertation we have successfully doped ZnO nanorods using Al as dopant using sonochemical method, showing an increased electrical conductivity, 10 times increase in carrier concentration and a transmission as high as 70%

4.1.1 Synthesis of Al-ZnO Nanorods

In order to grow Al:ZnO nanorod arrays, first a seed layer was deposited on n-type Silicon wafer as described earlier in chapter 2. Following the seeding process Al:ZnO nanorods were synthesized by sonicating a solution of 0.02 M ZNH, 0.02 M HMT and Aluminum nitrate (AlN) at 50% amplitude (~20 W cm-2) for 30 minutes. A solution of Ammonium hydroxide (NH₄OH) (28 wt% NH₃ in water, 99.9%, Aldrich) was added to maintain the pH of the solution around 6.2. Maintaining pH values around 6 is very crucial as sonochemical synthesis is highly sensitive to the changes in pH value of the solution. AlN has been used previously and proved to be a very effective 'Al' dopant in various hydrothermal methods. We have used two different concentrations of 0.0005 M AlN and 0.001M AIN to dope ZnO. An increase in concentration beyond 0.001 M did not result in nanorod growth. The structure and morphologies of the Al doped ZnO nanorods were studied using, FESEM (JOEL 7000), EDS (JOEL 6330) and X-ray diffraction (XRD). The transmission properties were studied using IPEC measurement system (PV Measurements) and Hall mobility and current voltage characteristics were measured using HMS 3000 equipped with 0.55 T permanent magnet and four point probe station equipped with HP 4455 semiconductor parameter analyzer.

4.1.2 Structural and Electrical Characterization

Figures 4.1 (a – c) shows the FESEM images of un-doped ZnO nanorods (S10) and Al doped ZnO nanorods with AlN concentration 0.0005 M (S11) and 0.001M (S12) respectively. It can be observed that vertically aligned ZnO nanorods were grown over n-Si substrate, with uniform density. The density of nanorods is slightly reduced when AlN was added, however there is not much reduction in the density between 0.0005 M AlN (S11) and 0.001 M AlN (S12). The EDS analysis is shown in figure 4.1 (d-e) indicates the presence of 'Al' in a very small quantities in both samples with 0.42 weight percentage and 1.28 weight percentage.



Figure 4.1 (a) SEM picture of undoped ZnO (S10) (b) ZnO nanorods doped with 0.5 mM AlN (S11) (c) ZnO nanorods doped with 1 mM AlN (S12) (d) EDS showing 0.42% weight percent of Al in S11 (e) EDS showing 1.23% weight percent of Al in S12

Figure 4.2 (a) shows the XRD analysis for samples S10, S11 and S12. The spectra were dominated by (002) peak of wurtzite ZnO, which implies that the preferred orientation is

along (002) direction. We did not observe any Al_2O_3 peaks as shown by other studies, but instead we observe the formation of double oxide gahnite (ZnO.Al₂O₃, JCPDS card No. 22-1034) similar to the one observed in ZnO nanowalls grown over 'Al' substrate. A similar peak was reported by Sining Yun *et. al* [14] and ascribed it to high aluminum concentrations. A broader (002) peak and peak shift towards higher angles of Al:ZnO nanorods indicate the lattice distortion due to the introduction of Aluminum.



Figure 4.2 (a) XRD analysis of S10, S11 and S12 samples showing a double gahnite peak and a peak shift in Al doped samples (b) Transmission spectra of undoped and Al doped ZnO (c) Current- voltage characteristics of S10, S11 and S12 after annealing at 300 C (d) Current- voltage characteristics of S10, S11 and S12 after annealing at 500 C after dopants are activated.

Figure 4.2 (b) shows the transmission spectra of S12 doped with 0.1 mM AlN and S10 an undoped ZnO nanorods. There is drop in transmission due to higher absorption in Al:ZnO however, albeit high transmission levels compared to other doped ZnO reported. Figure 4.2 (c-d) shows current-voltage characteristics of undoped and doped ZnO nanorods. It was observed that at 300 °C annealing temperature, the conductivity of undoped ZnO nanorods (S10) is high. This is expected as we have observed in ZnO nanowalls, the formation of double gahnite reduces the conductivity of ZnO nanorods. However, as the annealing temperature is increased to 500 °C the dopants are activated as such the conductivity of doped ZnO nanorods is higher than the undoped ZnO nanorods with S12> S11> S10. This behavior is collaborated by the Hall measurements which shows 10 times increase in carrier concentration when annealed at 500 °C.

4.2 P- Doping in ZnO Nanoshells

The role of oxide based functional materials has rapidly increased over the past few decades as the focus of research widened beyond their use as mere dielectrics. Today metal oxides find applications in catalysts, transparent electrodes, magnetic materials, bio sensors, piezo electric materials, photonics and many more.[15, 16] Among these oxide materials, Zinc Oxide (ZnO), a wide band gap (3.37 eV) semiconductor material with a high exciton binding energy (60 meV) at room temperature, has emerged as a promising material for fabrication of many optoelectronic devices such as light-emitting diodes (LEDs), laser diodes (LDs) and photodetectors (PDs).[17, 18] With a potential to replace other wide band gap materials such as GaN and SiC, ZnO has received substantial attention owing to its diverse morphologies, [19, 20] low toxicity, bio-compatibility, [21, 22] radiation hardness and amenability to wet chemical etching, which makes it appealing to commercialization and industry applications. One dimensional ZnO nanostructures such as nanorods, nanowires, nanotubes have been extensively studied in fabrication of optoelectronic devices due to their single crystalline nature and high surface to volume ratio. However, despite its many advantages, the progress of ZnO in nano-scale optoelectronic devices is hindered by the lack of stable and reproducible p-type doping. The difficulty in attaining p-type doping in ZnO has been attributed to various factors such as the self-compensating effect of the inherent intrinsic donor defects, hydrogen impurities and low solubility of ptype dopants. [23, 24] Apart from the above bottlenecks, the tendency of acceptors in ZnO to form a deep impurity levels instead of shallow acceptor levels has made it challenging to achieve p-type doping in ZnO. [25] This calls for development of a synthesis technique for successful and stable p-type doping in ZnO by introducing least amount of impurities without use of high processing temperatures and costs.

P-type doping in ZnO has been reported in recent times using techniques such as CVD, [26] ion implantation, [27] PLD, [28] MOCVD, [29] MBE, [30] which most commonly use nitrogen and phosphorous as dopant material. Aqueous solution methods have also been reported, which use both group I elements such as K, [31] and group II elements such as N, [32,33] and P, [34,35] but to the best of our knowledge there was no report on sonochemical method till date. In this research we for the first time synthesized p-type ZnO nanoshell (P:ZnO nanoshell) around ZnO nanorod (ZnO NR) by varying only the molar ratio of precursors and addition of ammonium dihydrogen phosphate (ADP) as

source of phosphorous (P) using sonochemical method. In comparison to the more conventional approaches such as hydrothermal method, sonochemical method is 10 times faster due to the fast hydrolysis rate caused by cavitation. Apart from that, sonochemical method is inexpensive, environmentally benign, CMOS compatible and has been proved to induce minimum defects in ZnO crystal unlike hydrothermal method. [36] Since sonochemical synthesis is carried out at atmospheric conditions, wide range of substrates such as flexible and wearable substrates can be used, which makes it very attractive for synthesizing ZnO nanostructures for various applications. Theoretical studies have shown that nitrogen (N) owing to its small atomic radii (1.71 Å) and low activation energy (0.32 eV) is the most promising dopant.[37] However, experimental results have shown that phosphorous can be successfully used as a p-type dopant, with measured ionization energies (0.18 eV) far less than theoretical calculations. The cause of this behavior is the formation of $P_{\text{zn}}-2V_{\text{zn}}$ complex defects, which act as shallow acceptors as studied by Woon- Jin Lee, Joongoo Kang and K. J. Chang and others. [38, 39] Ammonium dihydrogen phosphate (ADP), is the most suitable source of phosphorous for a solution based method as it readily dissolves in H_2O to form $Zn_3(PO_4)_2$ following the reaction:

$$[H_2PO_4]^- + [OH]^- \rightarrow H_2O + [PO_4]^{3-2}$$
$$2[PO_4]^{3-2} + 3Zn^{2+2} \rightarrow Zn_3(PO_4)_2$$

The $Zn_3(PO_4)_2$ contained in the ZnO nanorods decomposes at high annealing temperatures there by releasing P atoms, which would diffuse in to the ZnO crystal acting as acceptor dopants.[40]

Radial p-n homojunction offers wide range of advantages over conventional planar p-n junctions. Foremost, the p-n junction area is increased many fold in case of radial p-n junction when compared to planar p-n junction, this increases the electron-hole collection efficiency. The distance that the electron-hole pairs need to travel before they are effectively collected is very small when compared to planar p-n junction, which reduces the recombination losses encountered in optoelectronic devices. The radial p-n homojunction configuration also increases the light absorption, which helps in improving the efficiency of the devices. We used X-ray diffraction (XRD), Raman spectrum and current-voltage (I-V) measurements, which show a rectification, characteristic of a p-n junction to verify the formation of p-n junction. The capacitance – voltage measurements (C-V) were used to extract the doping concentration.

4.2.1 Growth of ZnO Nanorods

The ZnO nanorods were synthesized using the sonochemical method described in detail in chapter 3.[41] Briefly, 200 nm thick ZnO seed layer was initially deposited over an n-Si substrate by sonicating in a 100 mL solution containing 0.005 M Zinc acetate dehydrate (ZAD) in isopropyl alcohol (IPA) for four cycles of 15 minutes each. In the second step, high quality ZnO nanorods were grown by sonicating a 100 mL solution of 0.02M Zinc Nitrate Hexahydrate (ZNH) and 0.02M Hexamethylene tetramine (HMT) in deionized water (DI) for four cycles, each 30 minutes. The as grown ZnO nanorods are vertically oriented with an average diameter 50 nm and length 1 μm long.

4.2.2 Growth of P:ZnO Nanoshells

ZnO nanoshells were grown around vertically oriented ZnO nanorods by varying the molar ratios of ZNH and HMT rather than using any other surfactants. HMT is known to play a crucial role in promoting directional growth along the 'c' axis as it binds to the non-polar facets of ZnO, cutting off the supply of Zn^{2+} ions to these facets thereby promoting the growth in (0001) plane along 'c' direction.[42] Hence in this work, by changing the molar ratios of ZNH and HMT to 4:1 (1:1 is the preferred ratio for ZnO nanorod growth), we were able to promote the growth of ZnO along the non-polar facets of ZnO nanorods along with the growth in 'c' direction as described in chapter 3. This resulted in growth of a nanoshell structure enclosing the whole nanorod including top basal plane with increased junction area when compared to conventional core- shell nanostructures where the shell is grown only around the non- polar facets of ZnO nanorod. This method is advantageous over the more common ones, which use surfactants such as citrate ions to achieve directional growth, as no impurities were added in this method while forming a more continuous and defect free junction as observed from the TEM picture. [43-45] A continuous and defect free junction is desired in designing optoelectronic devices in order to minimize the recombination losses at the junction. P-type doping in ZnO nanoshells was achieved by adding Ammonium dihydrogen phosphate (ADP) to the precursors. The concentrations of ZNH and HMT were adjusted to keep the pH around 6, which is the optimum condition for sonochemical growth. For this purpose, 1.2 gm of ZNH, 0.14 gm of HMT and 0.0184 gm of ADP were dissolved in 100 mL of DI water and stirred for 5 minutes to form a homogeneous solution. The previously prepared sample with ZnO nanorods was dipped in to this solution and sonicated at 50% amplitude for 20 minutes. The thickness of the nanoshell depends on the sonication time. The thickness increases with increasing sonication time. As prepared samples were then cleaned in DI water and allowed dry in ambient.

The step by step synthesis of p-doped ZnO nanoshells around ZnO nanorods to form a homogeneous radial p-n junction is shown in figure 4.3. In step 1, ZnO nanorods were grown on an n doped Silicon (n-Si) substrate using the process described earlier. As synthesized nanorods are 1 µm long and 50 nm in diameter. In step 2, a 1 µm thick PMMA was deposited by spin coating followed by a hard-bake at 180 °C for 3 minutes. The PMMA is later carefully etched half way through, using reactive ion etching (RIE) in oxygen (O_2) plasma. This step exposes 600-800 nm of ZnO nanorods for further synthesis of p-doped ZnO nanoshells, while the remaining PMMA shields the n-type ZnO nanorods at their base. This prevents the growth of ZnO nanoshell around the core at the base, which would then make contact with seed layer shorting the device. In step 4, a 7 µm thick layer of AZ4620 photoresist was spin coated on the entire substrate and a 1 μ m × 1 μ m window was opened for the growth of ZnO nanoshells. In step 5, P:ZnO nanoshells were grown around the previously prepared ZnO nanorods by sonicating a solution of (ADP + ZNH + HMT) in deionized water (DI water) at 50% amplitude for 20 minutes. The concentration of ADP was varied from 1 mM to 2 mM in steps of 0.2 mM. Subsequently, the photoresist was removed by dipping in acetone for 5 minutes and cleaning with DI water. SEM analysis presented in figure 4.4 (a) and (b) have shown that after the growth of nanoshell, the effective diameter of the nanorods increases from 50 nm to 100 nm which indicates towards successful growth of nanoshell. In step 6, another layer of PMMA was deposited over the synthesized core-shell structures and etched in RIE to expose the tips of the nanoshells for depositing metal contacts for electrical measurements. However, due to a slight angular growth of some nanorods, not all nanoshells were exposed equally during this process and some of the nanoshells remained covered by PMMA as seen in figure 4.4 (d). Further etching of PMMA was not performed in order to avoid the risk of exposing ZnO core, which would result in the top metal contact making direct connection with n-ZnO core. In step 7, physical masks were used to deposit Ti/Au contacts on n-ZnO and Ni/Au on the P:ZnO nanoshells using e-beam evaporation. Current-voltage measurements at different annealing temperatures were taken, which show a successful p–type doping at 1.6 mM ADP and 500 °C annealing temperature.



Figure 4.3. Synthesis of p-doped ZnO nanoshells around ZnO nanorods: (a) Vertically oriented ZnO nanorods grown on n- Si; (b) A thin layer of PMMA spin coated to cover the ZnO nanorods; (c) Partially etched ZnO nanorods; (d) Growth of p-ZnO around partially exposed ZnO nanorods.

4.2.3 Deposition of Contacts

Electrical contacts on the ZnO nanorods were deposited using e-beam evaporation. Initially a 1 μ m thick PMMA was spin- coated on to the sample and was hard baked at 180 °C for 3 minutes. RIE was used to etch the PMMA to expose 100–200 nm of ZnO nanorods and P:ZnO nanoshells. Physical masks were prepared by using a laser engraver to cut windows with an area of 0.15 cm² in to a PMMA sheet. These physical masks were carefully laid on top of the already prepared samples and loaded in to CHA 10 keV four pocket e-beam deposition system and evacuated to 2 × 10⁻⁶ Torr vacuum. 30 nm/60 nm thick Ti/ Au and Ni/ Au contacts were deposited on the ZnO nanorods and P:ZnO nanorods and P:ZnO nanorods with an area of 0.15 cm² in to a PMMA sheet.

4.2.4 Structural and Electrical Characterization

Structural characterization was done by a 30 KV JEOL 7000 SEM and Philips CM-200 200 kV transmission electron microscope (TEM). Figure 4.4 shows the top view and cross-sectional view of SEM images of ZnO core and ZnO core-shell nanorods. It is observed that the ZnO nanorods, which act as core are 50 - 60 nm in diameter and are vertically oriented with a needle like pointed tips, while the newly formed nanoshell around the core is hexagonal in structure with a thickness of 40 - 50 nm giving an overall diameter of ≈ 100 nm. Figure 4.4 (c) shows the cross-sectional view of P:ZnO nanoshell/n-ZnO NR with PMMA filling the voids between the rods isolating the P:ZnO nanoshells from the

nanorods. Figure 4.4 (d) shows the exposed tips of P:ZnO nanoshell/n-ZnO NR for deposition of electrical contacts.



Figure 4.4. Structural characterization of ZnO core-shell by SEM: (a) ZnO nanorods with mean diameter of 50 nm; (b) ZnO core-shell structure with mean diameter 100 nm; (c) Cross-sectional view of P:ZnO nanoshell/ n- ZnO NR filled with PMMA; (d) partially etched PMMA and exposed tips of P:ZnO nanoshell for contact deposition.

Figure 4.5 (a-b) show the TEM image of a single core-shell nanorod and corresponding selected area electron diffraction pattern respectively. The ZnO core can be easily distinguished from the P:ZnO nanoshell due to an obvious difference in contrast between core and shell regions. From the SAED pattern, it is obvious that both the ZnO core and ZnO nanoshell that encloses ZnO rod are single crystalline and forms a defect free junction.

The result of the chemical composition analysis, which was performed using energy dispersive X-ray spectroscopy (EDS) is presented in Figure 4.5 (c), shows 0.96% weight of Phosphorous in the P:ZnO nanoshells. Figure 4.5 (d) shows the X-ray diffraction (XRD) measurements of both undoped ZnO nanorods and p doped P:ZnO nanoshell/n- ZnO NR structures after annealing. In ZnO nanorods, ZnO (002) diffraction peak is the dominant peak, which indicates that the samples have a preferred orientation along c-axis. We also observe ZnO (101) peak, which could be attributed to the tips of ZnO nanorods that sharpen like needles. In case of 'P' doped P:ZnO nanoshell/n- ZnO NR the (002) peak is the only dominant peak which corresponds well with the SEM pictures showing smooth hexagonal edges on the nanoshells. We also observe that compared with the ZnO NR, the position of (002) diffraction peak in P:ZnO nanoshell/ n- ZnO NR has shifted to the large-angle side by 1.2 degrees which provide a strong evidence that ZnO nanoshell has been successfully doped with phosphorous. Similar behavior was reported in earlier reports of 'P' doped ZnO. [40] In figure 4.5 (e) Raman spectra of the undoped ZnO nanorod and 'P' doped P:ZnO nanoshell are presented. The non-polar E_2 mode (E_2^{high}), which was observed at 436 cm⁻¹ in ZnO nanorods red shifts by 1.8 cm⁻¹ (to 434.2 cm⁻¹) and broadens in P:ZnO nanoshell. Such shift in the non- polar E₂ mode was reported earlier in Sb-doped and Ndoped ZnO. This decrease in frequency can be ascribed to the local lattice softening due to the incorporation of 'P' ions. Both XRD and Raman spectrum confirm the substitution of Zn sites by P ions in the crystal lattice and formation of $P_{zn} - 2V_{zn}$ complex defects. [46, 47]

Current-voltage (I-V) measurement is a useful tool to verify the formation of a p-n junction and to study the evolution of doping. Here, we studied the evolution of p-type doping by measuring the electrical current between the P:ZnO nanoshell and n-ZnO NR for various concentrations of ADP. For this purpose, Ni/Au contact was deposited on the top of P:ZnO nanoshells prepared with different concentrations of ADP and Ti/Au contact deposited on ZnO nanorods acts as ohmic contact on the n-side.



Figure 4.5. Growth orientation determination and chemical composition analysis by TEM, SAED, EDS, XRD and Raman: (a) TEM image of ZnO core-shell nanorod scratched from the substrate, shows a clear interface between nanoshell and core; (b) SAED pattern showing the nanorods are oriented along (002) plane; (c) EDS showing the presence of ZnO and P; (d) XRD of ZnO nanorods showing the growth dominated along (002) plane; (e) Raman spectrum of undoped and p-doped ZnO on ITO substrate where E_2^{high} peak is shifted by 1.8 cm⁻¹ towards the lower frequency.

Figure 4.6 (a) shows the schematic of P:ZnO nanoshell/n-ZnO NR p-n junction. Figure 4.6 (b-d) and (h) show the evolution of p-doping in ZnO as the concentration of ADP is increased from 1 mM to 1.6 mM. The rectification in current is observed when the

concentration of ADP is 1.6 mM shown in figure 4.6 (h), which proves successful p-type doping in ZnO nanoshells. The current – voltage measurement between two contacts placed on the same side have shown a linear behavior which confirms the formation of ohmic contacts and the rectification observed is entirely due to the junction between P:ZnO nanoshell and ZnO nanorod. Beyond 1.6 mM concentration of ADP the pH of the precursor solution becomes high making it untenable for growth using sonochemistry. It was also observed that temperature plays a crucial role in activating the dopants as can be seen from figure 4.6 (e-h), where I-V measurements of samples with same concentration of ADP and annealed at temperatures 25 °C, 300 °C, 400 °C and 500 °C were presented. We found that the dopants are activated after annealing at 500 °C for 30 minutes. This annealing temperature is required for the $Zn_3(PO_4)_2$ induced in to the ZnO nanorods during cavitation to decompose and release 'P' atoms, which would diffuse in to the ZnO crystal acting as acceptor dopants.[40] From figure 4.6 (h) we measured the junction properties. The turnon voltage is about 0.5 V and the reverse bias leakage current at -2 V is 200 μ A which increases with increasing reverse bias voltage. This could be due to non-uniform doping in the P:ZnO nanoshell region. The ideality factor of the diode 'n' can be calculated from:

$$n = \frac{q}{kT} \left(\frac{\partial \ln I}{\partial V}\right)^{-1} \tag{1}$$

where V is the applied voltage, k is the Boltzmann constant and the whole value of kT/q can be expressed as 0.0258 eV at T= 300 K. In the voltage range of 0.5 V and 2 V, the ideality factor 'n' was calculated to be 37. Such high ideality factors in wide bandgap semiconductors have been reported earlier. [48]



Figure 4.6. Verification of p-n junction by I-V measurements: (a) Schematic of P:ZnO nanoshell/ n- ZnO NR p-n junction; (b) I-V curve for 1 mM ADP doped ZnO nanoshell; (c) I-V curve for 1.2 mM ADP doped ZnO nanoshell; (d) I-V curve for 1.4 mM ADP doped ZnO nanoshell; (e) I-V curve for 1.6 mM ADP doped ZnO nanoshell at room temperature; (f) I-V curve for 1.6 mM ADP doped ZnO nanoshell at 300 °C; (g) I-V curve for 1.6 mM ADP doped ZnO nanoshell at 500 °C.

Capacitance voltage (C-V) measurements were measured using Keithley 4210- CVU to study the doping concentration of the as grown p-type ZnO nanoshells. When inverse of square of capacitance density was plotted against reverse bias voltage a linear plot is obtained. The slope of this plot is related to the doping density by:

$$N_{A} = \frac{-2}{q\varepsilon_{0}\varepsilon_{r}} \frac{1}{\frac{d(A/C)^{2}}{dv}}$$
(2)

where $q = 1.6 \times 10^{-19} \text{ C}$ is the charge of an electron, $\mathcal{E}_0 = 8.85 \times 10^{-14}$ is the permittivity of the free space, \mathcal{E}_r = relative permittivity of the medium, 'A' is the area of the contact and 'C' is the measured capacitance. Figure 4.7 (a, b) shows the capacitance vs voltage at

different frequencies and $(A/C)^2$ vs voltage of the p-n junction at 10 kHz frequency. The doping concentration N_A at 10 K Hz frequency was calculated to be 6.7×10^{17} cm⁻³ by measuring the slope of the C-V curve in the reverse bias presented in figure 4.7 (a).



Figure 4.7. Study of doping concentration by C-V measurements: (a) Capacitance vs Voltage measurements at frequencies between 5 k Hz to 10 k Hz; (b) $1/(C/A)^2$ vs Reverse Voltage at 10 K Hz frequency.

4.3 Conclusions

We have developed and optimized a process to successfully dope ZnO both n- type and ptype by using 'Al' and 'P' as respective dopants using sonochemistry for the first time. We have fabricated a radial p-n junction photo diode in a core-shell geometry. The core-shell structure was synthesized without using any external surfactants but rather used only different molar ratios of precursors. A doping concentration of 6.7×10^{17} cm⁻³ was achieved for the P:ZnO nanoshells. The as fabricated homogeneous p-n junction shows a rectifying behavior with an ideality factor of 37 and a reverse saturation current of 200 µA at -2 V. The developed method is fast, inexpensive and performed at ambient conditions while resulting very small amount of impurities. This facilitates the use of flexible substrates for

ZnO nanostructures, opening the door for flexible and wearable electronic and photonic

devices based on ZnO.

4.4 References

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CHAPTER 5

Zinc Oxide Based Biosensing

Zinc oxide nanostructures have been extensively used in fabrication of biosensors over the past few decades. The interest in ZnO nanostructures originates from their high biocompatibility, low toxicity and high isoelectric point. Zinc Oxide nanostructures due to high anisotropy can be grown in various morphologies and as such a myriad of nanostructures were studied for sensing various biological materials such as glucose, galactose cholesterol etc. During the course of this research we fabricated electrochemical immunosensors for detection of cortisol using ZnO NRs and ZnO NFs as immobilizing matrix and compared their performance and rate of electron transport.

5.1 Electrochemical Cortisol Immunosensor Based on Sonochemically Synthesized 1D ZnO Nanorods and 2D ZnO Nanoflakes

Electrochemical immunosensors demonstrate improved performance over other sensing methods with increased detection range, rapid selective detection, and higher sensitivity which make them the preferred sensing method of targeted biomolecules for health quality monitoring. In order to achieve low detection limits and high sensitivity the immobilizing matrix used in immunosensors should have high surface functionality, high biomolecule loading and small resistance to electron transport resulting in high electron transfer rate. [1,2] This encourages researchers to explore the utilization of inorganic semiconductor nanostructures (NSs) as potential immobilizing matrix for immunosensor development to detect desired target analyte. [2] ZnO, a semiconducting material with high catalytic

efficiency, biocompatibility, chemical stability in physiological environments, low toxicity and a high isoelectric point (IEP) of about 9.5, has all the qualities for biosensing applications.[3,4] Literature shows the use of ZnO-NSs with different morphologies such as nanorods (NRs), nanowalls (NWs), nanobelts and quantum dots as immobilizing matrix in fabrication of biosensors for the detection of physiologically relevant biomarkers such as cholesterol, galactose, glucose.[5-8] These sensors have demonstrated high stability, better sensitivity, fast response time and excellent selectivity. 1D ZnO-NRs which boast high surface-to-volume ratio, and high electron mobility [9] provides a direct, stable pathway for rapid electron transport and have emerged as a very promising material for immobilization of biomolecules without electron mediator for use in immunosensor. [10-12] However, despite high surface area in 1D ZnO-NRs most of the area is enclosed by non-polar(10-10) facets, which have less catalytic activity compared to (0001) plane. This results in reduced sensitivities due to diminished signal strength and weak adsorption of antibody. Apart from small signal strength, incompatibility with CMOS process, difficulty in making electrical contacts and conventional nanostructure synthesis processes that generally require high temperature hinder the integration of 1D ZnO-NSs in immunosensors.[14] 2D ZnO-NFs synthesized by sonochemical method have greater active surface area in (0001) plane. [19] This provides more binding sites for antibody in (0001) plane which has high surface charge density, higher catalytic activity, smaller charge transfer resistance and higher signal strength. Fabrication of electrical contacts on 2D nanostructures is much easier than 1D structures which allows for easier integration in to CMOS devices. Sonochemical process is a simple one-step process in which high quality single crystalline ZnO-NS's are grown at ambient conditions on a variety of substrates.

The method is rapid, inexpensive, low-temperature, catalyst-free, CMOS compatible, and environmentally benign. [14,15]

In this part of the research, a simple low-cost and label free immunosensors were fabricated using 1D ZnO-NRs and 2D ZnO-NFs as immobilizing matrix to detect cortisol. In brief, Cortisol is a steroid hormone known as a potential biomarker for psychological stress estimation. Abnormality in cortisol levels is an indicative of many other disorders such as Cushing's disease, inflammation and Addison's disease. [16, 17] Cortisol is an important molecule in regulating blood pressure, glucose levels, and carbohydrate metabolism. The conventional laboratory cortisol detection techniques such as chromatography, spectroscopy, immune-assays and ELISA are only limited to laboratory as they involve time taking measurement, complex sample preparation and expensive diagnosis. [16] This deemed the necessity of electrochemical immunosensing protocol for rapid and selective cortisol detection.

For selective cortisol estimation, Anti-cortisol antibody (Anti-C_{ab}) obtained from the BIO-MEMS & MICROSYSTEMS lab at FIU was immobilized on ZnO-NSs matrix via electrostatic interaction. The electron transfer phenomena for both ZnO-NSs during fabrication of immunosensor was studied using electrochemical impedance spectroscopy (EIS). The results of cyclic voltammetry (CV) studies related to electrochemical sensing showed a high sensitivity of 11.86μ A/M for sensor fabricated with ZnO NRs as immobilizing matrix and a low detection concentration of 1pM. For real applications, the developed immunosensor were tested on human salivary cortisol and results were validated using ELISA technique.

5.1.1 Synthesis of ZnO Nanostructures

Two different morphologies, 1D ZnO-NRs which are oriented perpendicular to the substrate and 2D ZnO-NFs were synthesized via sonochemical method described in previous chapter 3 and reported elsewhere. [18, 19]

For ZnO-NRs, 150nm gold was initially deposited on 'Si' substrate by e-beam evaporation followed by growth of ZnO seed layer and ZnO NRs. The seeding process was performed by immersing the substrate in a solution of isopropyl alcohol and 5 mM ZAD. The solution was then irradiated using a commercially available high-intensity ultrasound setup (750 W ultrasonic processor, Sonics and Systems), for 1-15 min cycles at an amplitude of 75% of the maximum amplitude (≈ 32.14 W cm⁻²) and frequency of 20 kHz. After the seed layer, nanorods were grown by immersing the substrate in an aqueous solution of 0.02M ZNH and 0.02M HMT and irradiating the solution at 50% of the maximum amplitude of the ultrasonic probe for 20 minutes providing the solution with an intensity of 21W.cm⁻². The substrate was rinsed with DI water then dried with N₂ after each cycle. For, ZnO-NFs, 150nm gold was initially deposited on 'Si' substrate by e-beam evaporation, followed by growth of ZnO-NFs. A solution of 20 mM ZNH and 20 mM HMT prepared at room temperature by stirring with a magnetic stir bar at 350 rpm for 5 min to ensure a mixed solution was used as the only solution for the growth of 2D ZnO-NFs. The solution with the immersed substrate was then irradiated at 70% of the maximum amplitude (≈ 30 W cm⁻²) and 20 kHz frequency for 15 minutes. This process was repeated four times to ensure a uniform growth. The temperature of the aqueous solutions did not exceed $\approx 70 \circ C$.

5.1.2 Fabrication of ZnO-NSs Based Immunosensors

For the preparation of ZnO-NSs based immunosensor, 10μ L of Anti-C_{ab} was covalently immobilized via electrostatic interaction. A huge difference in IEPs of ZnO (9.5) and antibody (4.5) ensures high binding strength using physical adsorption. Phosphate buffer saline (PBS, pH 7.4, 10 mM) solution was prepared by dissolving 1 PBS tablet in 200 mL of de-ionized water and was used to wash the ZnO-NSs/Au electrodes after they were incubated with Anti-C_{ab} for 2hrs to remove any unbound molecules. The non-binding sites of Anti-C_{ab}/ZnO-NSs/Au bio electrode were blocked by immobilizing 10µL of bovine serum albumin (BSA) for 1 hour followed by washing in PBS solution and drying at room temperature. The fabricated BSA/Anti-C_{ab}/ZnO-NSs/Au immunoelectrodes were stored in refrigerator at 4 °C when not in use.

5.1.3 Structural and Morphological Characterization

Scanning Electron Microscopy (SEM, JOEL 7000 FSEM) was used to image the surface during various stages of electrode fabrication. Figure 5.1 (a, b) show scanning electron microscopy (SEM) images of vertically oriented and uniformly grown ZnO-NRs and 2D ZnO-NFs over Au/Si substrate. The average diameter of ZnO-NRs is 20 nm and the average length is 3μ m. ZnO-NFs are 20nm thick and have the average lateral dimensions of 5μ m x 5μ m. Figure. 5.1 (a-b) shows the FESEM image of ZnO-NR/Au and ZnO-NF/Au electrodes before Anti-C_{ab} antibody was immobilized. Both anodes show a uniform and dense growth of ZnO nanorods and ZnO nanoflakes over gold coated silicon substrates.



Figure 5.1 (a) ZnO nanorods on Au/ Si substrate (b) ZnO nanoflakes grown in Au/ Si substrate.

Atomic force microscopy (AFM) is a powerful tool used to study the surface roughness variations at nanoscale. This helps us to understand and quantify the adsorption of antibody over the ZnO nanostructures after introduction of Anti-C_{ab}. The surface roughness of the electrodes prior to the introduction of antibody was measured at INSYST using atomic force microscope from Nanonics which is operated in resonance frequency mode using a 'Cr' tip with a resonant frequency of 34.95 kHz and tip diameter 20 nm. This roughness will be later compared with the one calculated after immobilization of antibody to study the amount of antibody that was immobilized. High Resolution Transmission Electron Microscopy (HRTEM), Tecnai transmission electron microscope with 300 keV electron beam was used to obtain, HRTEM and SAED data. The data confirms that both ZnO-NRs and ZnO-NFs are single crystalline and oriented in [0001] direction with wurtzite structure. Raman Spectroscopy was studied using Nomadic Raman microscope by BaySpec (532 nm laser).



Figure 5.2 (a) Illustration of ZnO-NRs and ZnO-NFs, prepared using sonochemical method along with immobilization of monoclonal anti-cortisol antibody onto ZnO-NSs to fabricate electrochemical cortisol immunosensor. (b) Electrochemical response at the fabricated electrodes.

In an electrochemical immunosensor, oxidation occurs at the biomolecule immobilized electrode-electrolyte interface, where the biomolecule is oxidized as it loses electrons which are collected by the ZnO nanostructures (Figure 5.2). The generated electrons are rapidly transferred through the electrode material to the external circuit. However, most biomolecules are insulating in nature which increases the charge transfer resistance inhibiting the charge carriers. This change in resistance is used in sensing the biomolecule otherwise called antigen as, clear and distinct oxidation peaks are detected for varying concentrations of biomolecule during CV measurements. Therefore, in fabricating electrodes for an effective biosensor, factors such as high antibody binding and electron transfer rate play a major role. 1D ZnO-NRs with high IEP (~9.5) has positive surface charge at pH value less than 9.5. All the experiments in this report were at a pH of 7.4
which was the recommended pH to presume higher biological activity. At this pH value 1D ZnO-NRs provides binding sites for antibody (Anti- C_{ab}) whose isoelectric point is 4.2 and thereby is negatively charged under the experimental conditions. The high surface area in 1D nanostructures when compared to bulk materials will provide much more binding sites thereby increasing the antibody loading. At the same time 2D ZnO-NFs also result in high antibody binding due to large area in (0001) plane which has higher catalytic activity and low charge transfer resistance.

EIS and CV studies were performed to compare the electro active behavior of the electrodes. The electro activity of ZnO-NRs/Au, ZnO-NFs/Au, Anti-C_{ab}/ZnO-NRs/Au and Anti-C_{ab}/ZnO-NFs/Au electrodes was studied using a three electrode system with 3 mL PBS (pH 7.4) containing 5 mM Fe(CN)₆^{3-/4-} acting as an electrolyte. VMP3 electrochemical workstation from BioLogic Science Instrument, France was purchased to perform the electrochemical analysis. All CV measurements were performed at 30 mV/s scan rate which ensured maximum and repeatable current values. The CV measurements were made between -0.6 V to 0.6 V multiple times in order to test the repeatability. The obtained results show that the ZnO-NSs are stable over a period of time. All the electrochemical measurements were made in triplet set and an average is used for analysis.

5.1.5 Electrochemical Sensing of Cortisol

Saliva samples were derived from a biologically-relevant sample of two female agricultural workers of reproductive age who engage in shift-work. Much evidence indicates that cortisol levels vary significantly among female shift workers, and that stress associated with agricultural work is associated with cortisol level variation, as is exposure to cortisol.

Farmworkers were 24 and 39 years of age, and were recruited based on the following criteria: be at least 18 years of age, work in agricultural fieldwork, work early morning shifts, and have exposure to pesticides. These samples were drawn from another study measuring the effects of a behavioral study to assess pesticide exposures and stress response. Recruitment, analyses and other procedures for this study were approved by the Human Subjects Boards at the Penn State University and the Florida International University.

Samples were collected five times: at awakening, before their first work break, at their lunch break (before eating), upon arrival from work and before bedtime. Samples were stored by participants at freezing temperature (-20°C) until they were transferred and stored at -80 by the project team in the field lab. The standard protocol requires saliva samples to be stored at -20 °C or lower for preservation of natural and biological properties over extended periods and while transportation. Each farmworker were given written and oral instruction on how to collect saliva using a Salivette (Sarstedt Inc., Rommelsdorf, Germany) that consisted of a small cotton roll inside a centrifuge tube. Participants collected a saliva sample by chewing on the cotton roll for 60 seconds.

The collected saliva samples were stored at -20 °C to maintain its biological characteristics. Before sensing, the saliva samples were de-freezed to room temperature and centrifuged at 3500 rpm for 15 minutes to extract saliva from salivettes. The separated saliva was then pipette out and kept at -20 °C when not in use. These samples were further used to detect cortisol using electrochemical immunosensors and ELISA for comparison.

The fabricated BSA/Anti- C_{ab} /ZnO-NP/Au immunoelectrodes were used to detect various concentration of cortisol using CV techniques in 30µL PBS (pH 7.4) containing 5

mM [Fe (CN)₆]^{3-/4-} as redox moieties. The fabricated immunoelectrodes were used to detect cortisol in saliva samples of two specimens collected at different time intervals. 10 μ l of saliva sample was placed on the electrochemical immunosensor and was allowed to incubate for 30 minutes to ensure proper binding. Immunoelectrodes were then washed using PBS ~10 mL to remove unbound saliva. Cortisol Enzyme Immunoassay Kit was procured from Arbor Assays, MI and a standard protocol was adopted to detect saliva cortisol. In Brief, 50 μ l of 4X diluted saliva was used for the detection of cortisol.

Figure 5.3 (a-b) shows SEM images of Anti-Cab/ZnO-NRs/Au and Anti-Cab/ZnO-NFs/Au electrodes, the change in morphology when compared to the SEM images acquired before introduction of antibody as shown in figure 5.1 (a-b) confirms the binding of antibody to nanostructure matrix. Figure 5.3 (c-d) shows the AFM images of the electrodes after introduction of the antibody. The average roughness of Anti-Cab/ZnO-NRs/Au (Figure 5.3 c) over an area of 20 µm x 20 µm was measured to be 124 nm while for the ZnO-NRs/Au electrode it was 171 nm. In figure 5.3 (d) the measured average surface roughness of Anti-C_{ab}/ZnO-NFs/Au was 158 nm, while ZnO-NF/Au electrode shows a higher average roughness of 214 nm. The drop in roughness value is due to immobilization of Anti-C_{ab} which fills the porous surface of the electrode. AFM measurements suggests that the ZnO-NFs structures offers higher surface area for successful a binding. The larger drop in surface roughness value in ZnO NFs after the introduction of Anti-Cab suggests greater antibody binding than was observed on ZnO NRs. This proves our initial assumption that higher surface area in (0001) plane observed in ZnO NFs can enhance the antibody absorption when compared to ZnO NRs whose directional growth results in smaller area in (0001) plane.



Figure5.3 (a) FSEM image of Anti-Cab/ZnO-NRs/Au immunoelectrode, (b) FSEM image of Anti-C/ZnO-NFs/Au immunoelectrode showing change in morphology after antibody introduced, (c) AFM image of Anti-C_{ab}/ZnO-NRs/Au with a measured roughness of 124 nm, (d) AFM image of Anti-C_{ab} /ZnO-NFs/Au immunoelectrode with a measured roughness of 158 nm, (e) Raman shift of Anti-C_{ab}/ZnO-NRs/Au immunoelectrode and (f) Raman shift for Anti-C_{ab}/ZnO-NFs/Au immunoelectrode.

Raman spectroscopy of Anti- $C_{ab}/ZnO-NRs/Au$ and Anti- $C_{ab}/ZnO-NFs/Au$ electrodes (Figure 5.3 e-f) shows many more peaks other than that characteristic of wurtzite ZnO-NSs, which suggests a successful binding of Anti- C_{ab} over ZnO-NSs surface. Further electro chemical impedance spectroscopy (EIS) is employed to study the electron transfer characteristics and association constant for binding interactions at the electrode surface. To improve the selectivity by locking the signal from unbounded sites on nanoparticles, 10µl of bovine saline albumin (BSA) solution is applied.

Figure 5.4 (a) shows the Nyquist plots of the EIS spectra for bare Au, ZnO-NRs/Au, Anti- C_{ab}/ZnO -NRs/Au, BSA/Anti- C_{ab}/ZnO -NRs/Au electrodes and the equivalent Randles circuit *(inset)*. The diameter of the semicircle represents the charge transfer resistance (R_{ct})

on the electrode surface. The value of R_{ct} changes from 108.8 Ω for bare Au electrode to 2.677 k Ω for ZnO-NRs/Au electrode, which further increases to 3.552 k Ω with the introduction of antibody. Upon introduction of BSA, an insulating material, it isolates the non-bounded sites and further increases Rct value for BSA/Anti-Cab/ZnO-NRs/Au electrode to 7.768 kΩ. The electron transfer from the electrolyte $(5mM [Fe (CN)_6]^{3-/4-}$ in phosphate buffer solution (PBS)) to the electrode is hindered with the addition of antibody which further decreases after addition of BSA, resulting in increase in R_{ct}. The same characteristics can be observed in BSA/Anti-Cab/ZnO-NFs/Au electrodes as shown in figure 5.4 (b). Here the R_{ct} values for bare gold electrode is 120.3 Ω , for ZnO-NFs/Au is 1.462 k Ω which increases to 1.643 k Ω upon application of antibody for Anti- $C_{ab}/ZnONFs/Au$ and finally to 5.962 k Ω upon introduction of BSA for BSA/Anti-C_{ab}/ZnONFs/Au electrode. As can be seen, the introduction of BSA, an insulating material greatly increases the charge transfer resistance which decreases the redox current. However BSA is essential in order to limit the binding of Cortisol at unwanted sites. We also observe that the charge transfer resistance of ZnO-NRs is higher than the electrodes with ZnO-NFs, which results in higher oxidation current at BSA/Anti-Cab/ZnO-NFs/Au electrode interface as more charge carriers are absorbed, when compared to BSA/Anti-Cab/ZnO-NRs/Au interface. This can be due to the high single crystalline nature of ZnO NFs when compared to the seed layer used in the growth of ZnO NRs. The thickness of ZnO NFs (≈ 200 nm) which is smaller than the seed layer (≈ 500 nm) used in ZnO nanorod growth can also play a major role in boosting the electron transfer from electrolyte to the electrode.



Figure. 5.4. Stepwise characterization related to the fabrication of ZnO-NSs based electrochemical immunosensor using electrochemical impedance spectroscopy and cyclic voltammetry (A and B). (a) An increase in R_{ct} from bare Au electrode to BSA/Anti- C_{ab}/ZnO -NRs/Au. (b) Rise of R_{ct} from bare Au electrode to BSA/Anti- C_{ab}/ZnO -NRs/Au immunoelectrode. (c) Fall in the magnitude of oxidation response current from Au electrode to BSA/Anti- C_{ab}/ZnO -NRs/Au immunoelectrode. (d) Fall in the magnitude of oxidation response current from Au electrode to BSA/Anti- C_{ab}/ZnO -NRs/Au immunoelectrode. (e) An increment in the magnitude of oxidation response current of Anti- C_{ab}/ZnO -NRs/Au immunoelectrode with increase in scan rate from 10 mV/s to 100 mV/s. (f) An increment in the magnitude of oxidation response current of Anti- C_{ab}/ZnO -NFs/Au with increase in scan rate from 10 mV/s to 100 mV/s.

The electron transfer rate (K_{et}) at the electrode interface can be estimated by using the EIS and equation (1) [21]

$$K_{et} = \frac{1}{2(R_{ct})(CPE)} \tag{1}$$

Where 'CPE' is constant phase element and R_{ct} is the charge transfer resistance. The binding constant K_{ac} which is an indicative of the stable binding between the nanoparticles and antibody can be estimated by using equation (2) [22]

$$K_{ac} = \frac{\Delta R_{ct}}{R_{ct}(C_o)} \tag{2}$$

Where ΔR_{ct} is the difference in charge transfer resistance due to the introduction of antibody and R_{ct} (c_o) is charge transfer resistance without antibody. The results of the EIS at various stages of electrode fabrication are given in Table 5.1. Anti-C_{ab}/ZnO-NFs/Au electrodes exhibited higher charge transfer rate when compared to Anti-C_{ab}/ZnO-NRs/Au which explains the higher oxidation current observed in CV measurements. As shown in figure 5.4 (c) and 5.4 (d), the magnitude of oxidation current reduces from bare Au electrode to BSA/Anti- C_{ab}/ZnO-NRs/Au and BSA/Anti-C_{ab}/ZnO-NFs/Au electrodes as the charge transfer resistance increases with every step. As indicated by the higher K_{ac} value in ZnO-NRs, the Anti-C_{ab}/ZnO-NRs/Au exhibit stable binding between Anti-C_{ab} and the NRs when compared to ZnO-NFs.

CV method was used to investigate the electro-active behavior of electrodes and determine the sensitivity of the fabricated electrodes over a range of Cortisol concentrations. Initially Anti- C_{ab} /ZnO-NRs/Au and Anti- C_{ab} /ZnO-NFs/Au electrodes were studied by varying the scan rate from 10 mV/s to 100 mV/s (Figure 5.4 e-f).

Electrode	$R_{ct}(\Omega)$	CPE1(µF/s)	K _{et}	Kac
Au	108.8	02.74	1677.22	
ZnO-NRs/Au	2672	09.80	19	
Anti-C _{ab} /ZnO- NRs/Au	3552	14.70	9.57	0.3268
BSA/Anti-C _{ab} /ZnO- NRs/Au	7768	10.45	6.159	
Au	120.3	10.11	411	
ZnO-NFs/Au	1459	20.83	16.45	
Anti-C _{ab} /ZnO- NFs/Au	1643	21.18	14.3	0.126
BSA/Anti-C _{ab} /ZnO- NFs/Au	7768	10.45	6.16	

Table 5.1. Summary of the impedimetric parameters obtained in stepwise fabrication of immunosensor

The change in current response with scan rate plotted in the inset of Fig. 4(e-f) shows that the oxidation current increases linearly with scan rate and follows the relation: (a) I (μ A) = 2.6 + 0.835v (mVs⁻¹) in case of Anti-C_{ab}/ZnO-NRs/Au electrode and (b) I (μ A) = 74.6 + 1.2v (mV s⁻¹) in case of Anti-C_{ab}/ZnO-NFs/Au electrode, which indicates a surface controlled process. 10 μ l of PBS (10mM, pH 7.0, 0.9% NaCl) solution containing specific concentration of Cortisol was introduced to BSA/ Anti-C_{ab}/ZnO-NRs/Au and BSA/Anti-C_{ab}/ZnO-NFs/Au electrodes by incubating for 30 mins. The prepared electrodes were cleaned using DI water before introducing Cortisol. During incubation, antigen (Cortisol) binds with the specific antibody (Anti-C_{ab}) forming an insulating immuno-complex between Anti-C_{ab} and cortisol which inhibits the electron transport from electrolyte to the electrode and further increases the charge transfer resistance (R_{ct}). Electrodes prepared from the same batch were incubated in solutions with different concentrations of Cortisol. The CV measurements indicate a fall in maximum oxidation current with the addition of cortisol to the electrode, confirming the increase in R_{ct}. This fall in value increases with the increase in concentration of Cortisol incubated in the range of 1 pM to 100 nM. The current difference directly correlates to the cortisol concentration being sensed. Sharp and distinct oxidation peaks were obtained for ZnO-NFs. For a concentration of 100pm, peak current of 100 μ A was observed for ZnO-NFs while for the same concentration the peak current was 50 µA for ZnO-NRs. High current values and well separated oxidation peaks were obtained for both Anti-Cab/ZnO-NRs/Au and Anti-Cab/ZnO-NFs/Au electrodes as shown in Figure 5.5 (a) and 5.5 (b). The graph showing change in current with the log of concentration is almost linear following the equation: (a) $\Delta I (\mu A) = 7.74 \text{ Log}$ (Cortisol conc. (M)) + 94.18 for BSA/Anti- C_{ab}/ZnO -NFs/Au giving a sensitivity of 7.74 μ A/M and (b) $\Delta I (\mu A) = 11.86 \text{ Log} (\text{Cortisol Conc. (M)}) + 121.69 \text{ for BSA/Anti-Cab/ZnO-NRs/Au}$ electrode at a sensitivity of 11.86 µA/M. A very clear separation of oxidation current is obtained, which shows that the electrodes are stable and can be re-used. EIS (Fig 5c & 5d) for both electrodes at different cortisol concentrations shows that the values of R_{ct} increases almost linearly with increase in cortisol concentration (shown in inset). CV based cortisol immunosensor were used to detect human salivary cortisol and obtained results were validated using commercially available ELISA sensing technique. The protocol adopted to detect saliva cortisol using electrochemical immunosensor and ELISA is reported in our previous work. [23]. Both ZnO-NSs based cortisol immunosensor are used for cortisol detection in human saliva collected at different time. The cortisol concentration of same

saliva samples is detected using ELISA and compared with the values measured using the fabricated immunosensor. It was observed that a factor of 2.6 was applied to all the CV calculations in comparison to ELISA.



Figure 5.5 (a) CV studies of immunosensors as function of cortisol concentration varying from 100 pM to 100 nM for BSA/Anti- C_{ab} /ZnO-NFs/Au electrode, inset: a calibration plotted between the obtained change in response current with respect to cortisol concentration (b) CV curves with cortisol concentration varying from 100 pM to 100 nM for BSA/Anti- C_{ab} /ZnO-NRs/Au immunoelectrode, inset: a calibration plotted between the obtained change in response current with respect to cortisol concentration (c) Change in EIS of BSA/Anti- C_{ab} /ZnO-NFs/Au immunoelectrode with respect to cortisol concentrations (d) Change in EIS of BSA/Anti- C_{ab} /ZnO-NRs/Au immunoelectrode with respect to cortisol concentrations

Salivary Cortisol (ELISA)(ng/dL)	Salivary Cortisol BSA/Anti-C _{ab} /ZnO-NFs/Au (ng/dL)	Salivary Cortisol BSA/Anti-C _{ab} /ZnO-NRs/Au (ng/dL)	
155	146	143	
144	139	138	
237	228	231	
72	65	62	
240	230	233	
153	142	141	

Table 5.2. Comparison of saliva cortisol estimated using ELISA and ZnO-NSs based cortisol immunosensor.

5.2 Conclusions

We have fabricated vertically standing 1D ZnO-NRs and lateral 2D ZnO-NFs on Au coated Si, using sonochemical technique to prepare label free, highly sensitive and selective electrochemical cortisol immunosensors. Both 1D ZnO-NRs and 2D ZnO-NFs exhibit selective electrochemical cortisol detection at 1 pM 100 times better than ELISA. ZnO-NFs with sharp and distinct oxidation peaks indicate a much better surface controlled process when compared to ZnO-NRs. The sensitivity of ZnO-NFs is 7.74 μ A/M which is comparable to 11.86 μ A/M of ZnO-NRs coupled with ease of fabrication in case of 2D materials, we have successfully demonstrated the use of 2D ZnO nanoflakes as an alternate material for use as immobilizing matrix in electrochemical biosensors. It is demonstrated in this work that ZnO-NSs based immunosensors are capable of detecting cortisol in human saliva sample and well-validated using ELISA values. This work suggest ZnO-NSs with high surface-to-volume ratio (NRs) and larger surface area (NFs) as viable materials for cortisol sensing in human saliva.

5.3 References

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CHAPTER 6

ZnO Nanostructures for Photovoltaic Devices

6.1 Dye Sensitized Solar Cells

In the 21st century, the world's demand for energy has reached unprecedented levels putting severe stress on world fossil fuel reserves. Fossil fuels have also been the main cause increased global carbon levels which has left us no choice but to develop renewable-energy technologies that can rival that of fossil fuels. Significant progress has been made in developing several technologies such as solar cells, fuel cells and bio fuels. Solar cells in particular have seen rapid rise in the past few years as can be seen in figure 6.1[1]



Figure 6.1. The rise of solar cell technologies over past 40 years

Conversion of solar energy in to electrical energy in solar cells is based on photovoltaic effect where solar energy is absorbed by photosensitive material to produce charge carriers which are collected by the electrical circuit. Several devices that make use of this technology have been developed over the past five decades as seen in figure 6.1 above. However, their use in daily life has been limited by (1) conversion efficiency and (2) cost per kW/h. The cost per kW/h has to reach well below US\$ 0.40 kW/h [2] before they can be commercially viable. Single crystalline Silicon solar cell, which was invented 50 years ago makes for 95% of the market share, [3] they operate on the principle of p-n junction formed at the interface of a p-type semiconductor and n-type semiconductor. In the presence of light, electrons and holes are photogenerated in the junction which are driven in opposite direction due to the internal electric field along the junction. These electron-hole pairs are collected by an external electrical circuit. Silicon (Si) solar cells can attain 92% of the theoretical energy conversion with maximum conversion efficiency reaching to 22% by Sun power. However, due to high material costs and low effective efficiencies 'Si' solar cells face difficulties in achieving wide spread usage. Several other technologies such as amorphous Silicon (a-Si), [4] polycrystalline solar cells [5] and compound semiconductor solar cells which use direct band gap semiconductors such as GaAs, CdTe, CIGS to achieve higher efficiencies [6-8] have been developed with the goal to achieve low cost per kW/h values. However, they still fail to replace the 'Si' solar cells as they have failed to achieve higher efficiencies and were plagued by low stability especially in a-Si solar cells, while compound semiconductors have high manufacturing costs.

The quest for new solar cell technologies that would realize the goal of attaining low cost per kW/h has received a significant boost during the past decade with the introduction of organic solar cells. Namely, two technologies (1) Dye sensitized solar cells (DSSC) and (2) Perovskite solar cells have shown a significant potential to reduce the operating and manufacturing costs while producing higher efficiencies. DSSC is a photoelectrochemical system in which dye molecules absorbed on the surface of porous metal-oxide film act as photo anode. A liquid electrolyte with Γ/I_3^- act as the medium between anode and platinum (Pt) coated fluorine doped tin oxide (FTO) glass which acts as cathode. When light is incident on DSSC the photons are absorbed by the dye monolayer creating excitons which are split at the dye and nanocrystalline metal oxide film interface. The electrons are fast absorbed in to the oxide film while the holes are released by the electrolyte and are collected by the cathode. The working mechanism is shown in figure 6.2 (b)



Figure 6.2 (a) ZnO nanorod based DSSC (b) Working mechanism of DSSC.

Compared to conventional silicon solar cells, the DSSCs offer an advantage in terms of manufacturing cost. As efficiencies reaching 10.4% in TiO₂- Ruthenium dyes [9] have been

achieved the scope of DSSCs replacing or being used in tandem with 'Si' solar cells has increased. This led to further investigation in to use of ZnO nanostructures in place of TiO₂. ZnO nanostructures provide significant advantage over TiO₂ due to their comparative high electron mobility and diverse morphologies which gives higher surface area than traditional TiO₂ nanoporous film. Table 6.1 shows the summary of DSSCs based on ZnO nanostructures reported in literature.

Structure	Photosensitizer	Efficiency (Max, Avg)
Nanoparticles	N719	2.1%, 1.5%
	N3	3.4%, 2%
Nanoporous films	N3	5.08%
	N719	4.1%,
	D149	4.27%
Nanowires	N3	4.7%, 2.4%
	N719	1.54%, 1%
Nanotubes	N719	2.3%, 1.5%
Nanobelts	N719	2.6%
Core-shell nanowires	N719	2.25%
(ZnO- TiO ₂)		
Aggregates	N3	5.4, 4%

 Table 6.1. Summary of DSSCs Base Don ZnO Nanostructures [10]

In this dissertation, we have fabricated DSSCs using ZnO nanorods Sonochemically grown over FTO and dipped in ruthenium dye (N3) as photo anode. The photovoltaic behavior of the fabricated DSSC was studied using Sol 2A solar simulator from Oriel. An efficiency of 2.74%, open circuit voltage of $V_{oc} = 0.843$ V, short circuit current (I_{sc}) of 6.46 mA and high fill factor of 50.3 were measured. These values are in the range of best efficiencies achieved using ZnO nanorods as shown in Table 6.1

6.1.1 Fabrication of Dye Sensitized Solar Cells

The fabrication of ZnO DSSC starts with the synthesis of ZnO nanorods on FTO. This is done by the process discussed in chapter 3. As synthesized ZnO anodes are shown in figure 6.3 (a)



Figure 6.3 (a) ZnO nanorods grown over FTO (b) Photovoltaic behavior

The dye which acts as photo absorber was prepared by dissolving 17.8 mg of N3 dye in 50 ml of ethanol to make 0.3 mM solution. The immersion time was optimized by immersing ZnO/FTO samples in dye solution for various times. It was found that 3 hours immersion time produces the best results. Beyond 6 hours of immersion, the dye molecules begin to agglomerate around the ZnO nanorods which decrease the charge transfer efficiency of ZnO nanorods. After immersing the ZnO/FTO samples in dye for 3 hours, the anodes were removed and washed in ethanol to remove any excess dye. Cathodes were prepared by spin coating a solution 'Pt' nanoparticles (platisol obtained from Solaronix) and sintered at 400 °C for 30 minutes. The cathode and anode were brought together and sealed using Meltonix hot melt sealing film from Solaronix. Iodolyte HI-30 a 30 mM Iodine/tri-iodide electrolyte

procured from Solaronix was introduced through the holes drilled on the back surface of the cathode which were then sealed using epoxy from Solaronix (Amisol).

6.1.2 Optical & Electrical Characterization

The fabricated DSSCs were characterized using Oriel Sol 2A solar simulator under 1 Sun shown in figure 6.4 (a). The resulting photovoltaic behavior is shown in figure 6.3 (b). The DSSCs presented an efficiency of efficiency of 2.74% which is in accordance with the average efficiencies obtained using ZnO nanowires as shown in table 6.1. However the fabricated DSSCs suffered from well-known limitations that have plagued ZnO based DSSCs as, the efficiencies fell rapidly due to rapid degradation of ZnO nanorods in acidic N3 dye and the gradual leakage of liquid iodide electrolyte. Apart from the above limitations we have also observed that the thick seed layer used in sonochemical synthesis of ZnO nanostructures has led to high recombination losses resulting in low short circuit current (J_{sc}). This led us to consider improved and stable organic solar cells that have inspired research in photovoltaics recently namely, Perovskite Solar cells.



Figure 6.4 (a) Solar Simulator ORIEL Sol 2A (b) Spectral Response Measurement System

6.2 Perovskite Solar Cells

The difficulties in fabricating stable DSSCs has inspired investigators to explore photon absorbing materials other than ruthenium dyes. In year 2006, Miyasaka and co-workers reported the use of CH₃NH₃PbBr₃, a perovskite material as a sensitizer in photoelectrochemical cell and reported an efficiency of 2.2%. [11] One year later they reported an improved efficiency by replacing Bromine with Iodine CH₃NH₃PbI₃. [12] However, these cells suffered from the same issues that plagued DSSCs as the perovskite easily dissolves in the liquid electrolyte. This problem was later overcome by replacing the electrolyte with solid hole-transportation layers. This led to higher efficiencies and increased stability. [13] There were even reports where no hole-transportation layer was used and instead, the perovskite layer acts as both photo sensitizer and hole conduction layer. [14, 15] Recent efforts have resulted in efficiencies greater than 15% [16] with high open-circuit voltage V_{oc} and high short circuit currents (J_{sc)} owing to high diffusion lengths in perovskite. PCE of 20% is predicted from perovskite solar cells.[17] This indicates that organometal halide perovskite is a very promising material for future solar cells.

There are two configurations in which perovskite solar cells are being developed. One is based on electron injection in to an oxide layer such as TiO_2 and ZnO. The other is based on scaffold layer or p-n junction type planar structure. One dimensional ZnO nanostructures are attractive alternative to TiO_2 as they provide faster electron mobilities, however the photovoltaic performances of ZnO nanorod based perovskite solar cells has been poor when compared with TiO_2 . This has changed when Dae-Yong and his colleagues reported 11% efficiency in ZnO nanorod based perovskite solar cells. [18] This has inspired us to fabricate ZnO nanorod based perovskite solar cells. For fabrication of perovskites owing to the lack of stability of CH₃NH₃PbI₃ in conditions where humidity is higher than 37% and also keeping in mind the toxicity of Pb, it is imperative to have working glove box for fabrication of perovskite solar cells. As such we have installed mbraun's labstar glove box with nitrogen as working gas shown in figure 6.5. A compact spinner and a hotplate were placed inside for fabrication of perovskite. A small desiccator with a hand pump for creating vacuum was used for sample transportation. The humidity and oxygen content in this glove box is maintained at 0.1 ppm. The fabrication of the cell was based on the findings of Dae-Yong son and colleagues.



Figure 6.5 MBRAUN glove box filled with nitrogen with O₂ and H₂O maintained less than 1 part per million

6.2.1 Fabrication of Perovskite:

FTO glass substrates were cleaned with acetone, isopropyl alcohol followed by DI water several times. The cleaned substrates were used as substrates for growth of 1 μm long ZnO nanorods by sonochemical method as described in chapter 3 of this dissertation. Figure 6.6 (b-c) shows both the top view and cross-sectional view of ZnO nanorods grown over FTO.



Figure 6.6 (a) Schematic of fabricated Perovskite solar cell (b) Top view of ZnO nanorods on FTO (c) Cross-sectional view of ZnO nanorods on FTO (d) Perovskite filling on ZnO nanorods/FTO (e) Cross-sectional view of perovskite/ZnO nanorods/FTO (f) Perovskite filling on ZnO nanorods/FTO at low spin speed

A two-step method was used to deposit perovskite on ZnO nanorods. [19] A 1 M PbI₂ solution was prepared by dissolving 4.6 g of PbI₂ in 10 ml of N, N- dimethylformamide (DMF) and stirred on spin coater while temperature is maintained at 70 °C. This solution is allowed to stir at the set temperature for 2 hours. 100 μ l of this solution was dropped on ZnO NR/FTO substrate and spun immediately at 3000 rpm for 20 seconds and immediately dried at 100 °C for 10 minutes. After the substrate is cooled, a 200 μ l of CH₃NH₃I was

dropped and allowed to rest for 20 seconds, before it was spun at 2000 rpm for 20 seconds. The sample was placed on hot plate at 100 °C and dried for 10 minutes. As the sample dries, the layer on the ZnO nanorods turns brown in color from yellow, indicating the formation of perovskite. After the samples were cooled, a physical mask was placed on the samples and transferred in to a compact desiccator for transferring in to e-beam deposition chamber for gold contact deposition. An extra sample was prepared for SEM and XRD analysis. Photo- current vs voltage characteristics were measured using Oriel SOL 2A solar simulator.

6.2.2 Structural & Electrical Properties

SEM analysis was done using JEOL 7000 FESEM to study the deposition of perovskite on ZnO nanorods. As can be seen in figure 6.6 (c-d) the perovskite failed to cover the whole ZnO nanorods as tips of nanorods are clearly visible protruding from the perovskite layer. This resulted in shorting the gold contacts to FTO as we observed no photovoltaic activity. In second phase, the thickness of the perovskite was increased so as to cover the as synthesized ZnO nanorods by reducing the spin speed of PbI₂ to 1000 rpm from 3000 rpm. Figure 6.6 (e) shows SEM picture of Perovskite/ZnO nanorods/FTO when perovskite was deposited at lower spin speeds. It can be observed that at lower speeds the perovskite film is non uniform however no ZnO nanorods are visible. Figure 6.7 (a-c) shows the XRD analysis which shows the formation of perovskite on ZnO nanorods. Figure 6.7 (a) shows the transition from lead iodide to methyl ammonium lead iodide in samples S1 as PbI₂ peaks disappear and perovskite peaks appear. The perovskite peaks indicate towards formation of tetragonal crystal structure in perovskite.



Figure 6.7 (a) XRD showing the change in peaks as PbI_2 is changed to perovskite after introduction of CH_3NH_3I and heating. (b) Perovskite formed in cubic lattice at 3000 rpm spin speed (c) perovskite formation in cubic lattice at 1000 rpm spin speed and showing less peaks of ZnO due to higher coverage.

Figure 6.7 (b) shows the XRD analysis of sample S2 in which the spin speed was maintained at 3000 rpm resulting in incomplete coverage of ZnO nanorods. The dominant (100) clearly indicates towards the formation of cubic lattice. The difference in crystal structure when compared to sample S1 could be due to experimental conditions as sample S1 was prepared in clean room for experimental convenience which could have resulted in tetragonal crystal structure due higher humidity present. Figure 6.7 (c) shows the XRD analysis on Sample S3 on which PbI₂ was deposited by spin coating at 1000 rpm speed which led to complete coverage of ZnO nanorods albeit the non-uniformity. We observe a clear dominant (100) peak indicating a cubic lattice structure. The lack of intermediate peaks corresponding ZnO unlike the previous samples also indicate towards complete coverage of perovskite. However the sample S3 when characterized under solar simulator again showed very small short circuit currents J_{sc} of 54 μ A which leads to very low efficiencies. The extremely low values of short circuit currents indicate towards very high recombination losses which can attributed to the thickness of ZnO seed layer which is very high. However from our previous experiments we realize that a thick seed layer ≈ 400 nm

is essential for uniform and vertically oriented ZnO nanorods. This calls for alternative methods to grow ZnO nanorods.

In this work, we used ZnO nanoflakes/ graphene as a seed layer for ZnO nanorod growth. As shown in chapter 3, the ZnO nanoflakes are extremely thin with maximum thickness reaching 200 nm. ZnO nanoflakes which are oriented in (001) plane are single crystalline in nature and form a thin continuous layer on graphene. In figure 6.8 (a) we show the modified architecture of ZnO nanorod based perovskite using flexible PET as substrate. Use of ZnO nanoflakes not only reduces the recombination losses but also enhances the electron transport as was shown in earlier studies (cortisol sensor). Figure 6.8 (b) shows the successfully grown ZnO nanorods using ZnO nanoflakes grown over graphene/PET as seed layer. Further growth cycles are needed to achieve a uniform vertically oriented ZnO nanorods. Another architecture we use is ZnO Nanowalls grown on 'Al' coated PET substrate as scaffold layer in fabricating planar perovskite solar cells as shown in figure 6.8 (c). ZnO Nanowalls are 10- 15 nm thick and highly porous which is a required feature for successful fabrication of perovskite solar cells.



Figure 6.8 (a) ZnO nanorods grown over ZnO nanoflakes/graphene/PET as electron transport layer (b) Partially grown ZnO nanorods on ZnO nanoflakes/PET showing both ZnO nanorods and graphene (c) ZnO Nanowalls grown over Al/PET as scaffolding layer planar perovskite solar cell.

6.3 Conclusions

In conclusion, we have fabricated DSSC's using ZnO nanorods as electron transport material. The as fabricated DSSC' showed an external quantum efficiency of 2.74% which is in accordance with the values reported in literature. We also observed that ZnO is not stable in acidic dyes which prompted us to fabricate perovskite solar cells. We have fabricated perovskite solar cells using ZnO nanorods grown on FTO as electron transport layer and found that the thick seed layer used in sonochemistry results in heavy recombination losses. To overcome this we have come with a new method of synthesizing ZnO nanorods, where ZnO nanoflakes grown over graphene/PET are used as electron transport layer in fabrication of flexible perovskite solar cells.

6.4 References

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CHAPTER 7

Photodetectors

7.1 P:ZnO-ZnO Core-Shell Photodetectors

Zinc oxide due to its wide bandgap and superior stability under harsh environments has been widely used for UV photodetection.[1] In a ZnO p-n junction, an increase in reverse current is observed under UV illumination as electron-hole pairs are generated in depletion region by band to band excitation in semiconductor when $hv > E_g$. These electron-hole pairs are quickly swept away towards the electrodes due to the internal electric field. This increased reverse saturation current due to the incident light is also called the photo current (I_{ph}).[2] We have fabricated a UV p-n junction photodetector by growing a P:ZnO nanoshell/ n- ZnO NR p-n junction on Au/Ti/Si substrate which acts as ohmic contact to the ZnO nanorods while a physical mask with a window size of 0.15 cm² was used to deposit 100 nm thick ITO by RF sputtering to form contact on the P:ZnO nanoshells to allow maximum transmission of incident light. The RF sputtering was performed at room temperature with pressure maintained steady at 11 mTorr. The RF power and time were 70 W and 100 minutes respectively to produce a deposition rate of 1 nm/s. The schematic of as prepared p-n junction is shown in figure 7.1 (a). ZnO with wide bandgap (3.37 eV) is known to generate electron-hole pairs when exposed to UV illumination where wavelength corresponds to the band edge excitation. The diode was tested under a commercially available UV compact fluorescent lamp with a broad band emission (365 nm - 385 nm). The output of the lamp measured by OAI 306 UV power meter was 4.6 mW/cm². The photoresponse under UV illumination is measured and plotted in figure 7.1 (b), which

shows a 100 fold increase in reverse bias current at 4 V reverse bias, which gives a photo current (I_{ph}) of 371.88 µA. The photo responsivity R_{ph} measures the input-output gain of the photo detector given by:

$$R_{ph} = \frac{J_{ph}}{P_o}$$
(3)

Where J_{ph} is the photo current density, and P_0 is the input optical power density. We have estimated the number of P:ZnO nanoshell/n-ZnO NR structures that actually make contact with our deposited ITO contact and calculated the area under the contacts from which J_{ph} and R_{ph} were calculated to be 0.02 A/ cm² and 60 A/W respectively.



Figure 7.1. Characterization of the fabricated UV radial p-n junction photo detector: (a) Schematic of p-ZnO/n-ZnO core-shell homojunction; (b) Photoresponse in presence of UV light.

Another important feature of photodetectors is their noise equivalent power (NEP), which represents the lowest input power that gives a unit signal to noise ratio. Conventionally, in p-n junction photodiodes the contribution of shot noise towards over all noise is negligible. However, considering the 200 μ A reverse saturation current in our diode, we have considered both shot noise and Johnson noise for noise calculations. Johnson noise current

 i_j is due to thermal noise associated with the resistance and is given by $i_j = \sqrt{\frac{4kBT}{R_{eq}}}$, where

k is the Boltzmann constant, T is the absolute temperature and R_{eq} is the equivalent resistance. The R_{eq} is obtained from the slope of I-V of P:ZnO nanoshell/ZnO at -2 V measured in dark. For R_{eq} of 4.45 k Ω and for normalized bandwidth (B) at room temperature the Johnson noise i_j is 19.2 μ A. The shot noise is given by $i_s = \sqrt{2q(I_{ph} + I_d)B}$, where I_{ph} is photo current and I_d is dark current. The measured shot noise current at -2V is 7.28 pA. The total noise current $i_n \approx i_j = 19.2 \ \mu$ A and the NEP which is the ratio of i_n and the responsivity is calculated.[3, 4] We obtain a minimum NEP of 324 nW/ \sqrt{Hz} . Table 7.1 compares the performance of as fabricated P:ZnO nanoshell-ZnO p-n junction photo diode with other UV photo detectors reported in literature.

Material/	Device	Method	Responsivity	Turn On
Structure				Voltage
P- NiO/ n- ZnO	Planar p-n	RF Sputtering	0.09 A/W	0.9 V,[5]
	junction diode			
Graphene/ n-	Schottky		113 A/W	0.5 V[6]
ZnO	junction diode			
P- GaN/ n- ZnO	Planar p-n	MBE	0.68 mA/W	3.7 V[7]
	junction diode			
P (Al-N) ZnO/ n-	Planar p-n	Sol- Gel	4 A/W	1.4 V,[8]
ZnO	junction diode			
P (Sb) ZnO/ n-	Radial p-n	Hydrothermal	0.425 A/W	2.5 V,[4-26]
ZnO	junction diode	method		
P-ZnO/ n-ZnO	Radial p-n	Sonochemical	60 A/W	0.5 V [present
	junction diode	Method		work]

Table 7.1. Performance comparison of the fabricated UV photo detector with the recently reported ones

7.2 Highly Sensitive Wide Bandwidth Photodetector Based on Internal Photoemission in CVD Grown P-Type MoS₂/Graphene Schottky Junction

Graphene, a semi-metallic, two-dimensional (2D) material with a single atomic thickness has been the most widely researched material owing to its superior electrical, mechanical and optical properties such as high carrier mobility ($\approx 200, 000 \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$), high mechanical strength, flexibility, transparency, and high absorption rate (2.30% of incident light per each layer) among many other properties.[10, 11] Further, gapless nature of graphene enables charge carrier generation over a wide energy spectrum from ultra violet (UV) to terahertz (THz) making it a unique material for photonics and optoelectronics during the recent past.[12] Making use of the unique capabilities of graphene, many studies have been performed over the past decade to demonstrate highly efficient graphene photodetectors (GPDs) from visible to infrared (IR) range.[13-15] Despite the many favorable properties, factors such as low absorption due to single layer nature of graphene limited the responsivity of conventional GPDs to few tens of mAW⁻¹.[12] Alternative methods such as colloidal quantum dots, microcavities or plasmonic nanostructures have been used to improve the responsivity of the GPDs but these devices were found to respond to very limited wavelengths of light.[16] On the other hand, 2D semiconductor transition metal dichalcogenides (sTMDs) materials have also been widely used for photodetector applications.[17] Atomically thin sTMDs provide a wide range of advantages in optoelectronics when compared to their bulk counterparts owing to their high transparency, flexibility and tunable bandgap. In particular, MoS₂ is a unique semiconductor material whose bulk counterpart has an indirect-gap of 1.20 eV but when reduced to a monolayer, it changes into direct bandgap material with a bandgap of 1.80 eV due to quantum confinement. This results in higher absorption coefficient and efficient electron-hole pair generation under photo excitation in monolayer MoS₂. A 10⁴ fold enhancement of luminescence quantum yield for monolayer MoS₂ compared with bulk MoS₂ has been observed. [18] There are several reports on the fabrication of MoS₂ based photodetectors. The first phototransistor based on monolayer MoS₂ showed a relatively poor photo responsivity of 7.50 mAW⁻¹.[19] Woong Choi *et al.* used multilayer MoS₂ to achieve 100 mAW⁻¹ photo responsivity,[20] showing an improvement over the first phototransistor but still low compared to GPDs. Since then, much progress has been made and more recently, phototransistors with responsivities ranging from few hundreds of mAW⁻¹ to hundreds of AW⁻¹ were reported. [21, 22] However, the reported MoS₂ based photodetectors are limited by narrow wavelength range and long response times.

To further improve the properties of graphene and MoS₂-based photodetectors, it is imperative to design new device structures like composites, heterostructures or functionally graded thin films. Numerous efforts have been made to combine the properties of TMDs and graphene by assembling graphene with other 2D crystals to create multifunctional high-performance hybrid devices. [23- 29] One of the hybrid structures includes the combination of graphene and MoS₂ in the form of thin film heterostructures. Coupling graphene with MoS₂ can produce a hybrid material that can utilize the high photon absorption capability in MoS₂ and high electron mobility in graphene to realize a highly efficient photodetector as shown by Wenjing Zhang *et al* who have achieved a high responsivity of 1.2×10^7 A/W by applying a gate voltage of -10 V at 690 nm laser illumination by simply stacking n-type monolayer MoS₂ over graphene.[30] In this research, we have assembled few layers of graphene with few layers of *p*-type MoS₂ to form a simple MoS_2 /graphene Schottky junction photodiode. We found that *p*-type MoS_2 provides a much lower Schottky barrier height with graphene with a broad spectral range of detection.

7.2.1 Fabrication of MoS₂/Graphene Photodetector

Fabrication of MoS₂/graphene junction devices starts with the growth of atomically thin MoS₂ film on Si/SiO₂ substrates followed by the transfer of graphene. The growth of both MoS₂ and graphene was done by "Nanomaterials and Device Laboratory" at University of North Texas. For MoS₂ deposition, the first step involves the sputtering of Mo thin films on (100) oriented *n*-type (As doped, resistivity $< 0.005 \ \Omega.cm$) silicon substrates coated with 300 nm thick SiO₂ layer. High purity (99.99%) Mo metal target of 50 mm diameter was used for sputtering Mo thin films. In second step, magnetron sputtered Mo/SiO₂/Si films were subsequently placed in a low-pressure chemical vapor deposition (LPCVD) system (Graphene Square CVD) and sulfurization was proceeded at 600 °C for 30 min. [31] Polycrystalline copper foils (Nimrod Hall, 99.9% purity, 25µm thick) was used as the substrates for graphene growth. The 'Cu' foil was electro-polished using a home-built electrochemistry cell followed by rinsing in methanol and deionized water solution (1:1), and then gently blown with compressed nitrogen for drying. The cleaned 'Cu' foil was loaded into the low-pressure thermal CVD system. The reaction chamber was evacuated to ~0.1 mTorr and refilled with 2.5 sccm of pure H_2 . The temperature was increased to 1000 ^oC within 40 min and then annealed for 1 h. After annealing, 20 sccm of CH₄ and 2.5 sccm of H₂ were introduced into the chamber for graphene growth at 1000 ^oC under the pressure of 5 Torr for a certain growth time. After the growth, the CH₄ flow was turned off and the

chamber was cooled down under H₂ environment. [32] The graphene transfer method used was as follows: (1) PMMA (MicroChem Corp. 495 PMMA A2) solution was spin-coated onto the graphene/Cu foil at 5000 rpm for 1 min, followed by drying the sample in the air; (2) the Cu foil was etched with metal etchant (type I, Transene company, INC.), resulting in the PMMA/graphene film floating to the top of the etchant solution; (3) the PMMA/graphene film was washed in a beaker containing deionized water, and floated on 10 % HCl solution for 10 min, and transferred to a beaker with deionized water for another wash (done $3\times$); (4)the film was transferred to a beaker with deionized water for another wash and this process was repeated three times; (5) the film was then finally scooped onto a Si/SiO2/MoS₂ substrate with the PMMA side up and dried in the air; (6) the PMMA layer was removed by immersing the sample in a large volume of acetone; (7) the sample was dried with compressed nitrogen. Figure 7.2 (a) shows the schematic of the $MoS_2/graphene$ photodetector device structure simply formed by a manual stacking of a graphene onto MoS₂. 50 nm thick Au/Ti metal contacts were deposited by electron beam evaporation after using a shadow mask. The shadow masks were carefully prepared using a laser writer so as to maintain the contact area and position accurately. The deposited contacts were later annealed at 350 °C for 30 minutes.

7.2.2 Structural Properties of MoS₂/Graphene

Figure 7.2 (a) shows the schematic of the studied photodiode. Figure 7.2 (b) shows the optical microscope image of MoS_2 /graphene junction which clearly shows a color contrast over the four different regions; (1) SiO_2 (2) MoS_2 (3) Graphene and (4) MoS_2 /graphene junction (MGJ). Atomic force microscopy (AFM) was used to study the thickness and the

interface of MoS_2 and graphene. Figures 7.2 (c) and 7.2 (e) show the thickness measured using AFM. We observed that the average thickness of MoS_2 layer is 2 nm and that of graphene is 5 nm, which suggests that as grown MoS_2 was about 3-4 layers thick. Raman spectroscopy which is a widely used method to determine the number of layers, quality and defects present in these 2D materials [33] was used to complement the AFM findings and to study the quality of MoS_2 and graphene films.



Figure 7.2. (a) Schematic of p-type MoS_2 /graphene junction device (b) Optical image of MoS_2 /graphene junction with individual areas clearly visible (c) AFM height profile showing height of MoS_2 (d) AFM image of MoS_2 -SiO₂ interface (e) Thickness of graphene at MoS_2 /graphene junction (f) AFM image of MoS_2 /graphene junction
Figure 7.3 (a) shows the Raman spectra of MoS_2 and graphene, measured over SiO_2 and junction regions using a 532 nm excitation laser line.



Figure 7.3. (a) Raman spectrum of MoS_2 and Graphene on SiO_2 and junction (b) Photoluminescence of MoS_2 on SiO_2 showing sharp peaks at 1.7, 1.75 and 1.85 eV, and smaller ones at 1.4 and 1.55eV.

We observed strong peaks at 379.77 cm⁻¹ and 403.52 cm⁻¹ corresponding to the in plane E_{2g}^{1} mode and out-of-plane A_{1g} mode of MoS₂, respectively. The peak frequency difference of 23.75 cm⁻¹ between the two modes indicates that as grown MoS₂ is 4-5 layers thick, which is in agreement with the previously reported results [34] and also with our results obtained by AFM height analysis. Raman spectra of graphene on SiO₂ shows the formation of '2D' and 'G' characteristics peaks at 2666.83 cm⁻¹ and 1580.16 cm⁻¹, respectively, which is consistent with the previous reports on few layer graphene.[35] However, a red peak shift of ~18.8 cm⁻¹ was observed in the '2D' peak of graphene when Raman spectra was taken at the junction. The observed peak shift could be attributed to an increase in hole concentration at the MoS₂/graphene junction. [36] This is a common phenomenon for *p*-

type semiconductor-metal Schottky junction where holes travel from the p-type semiconductor to metal as their Fermi levels are aligned. Figure 7.3 (b) shows photoluminescence (PL) spectrum measured at room temperature.

Peak 'A' associated with excitonic transitions at the 'K' point of Brillouin zone is observed at 1.85 eV while peak 'I' at 1.40 eV corresponding to indirect gap which is prominent in bulk MoS₂ is almost negligible confirming that the fabricated device is formed by few layers of MoS₂. [37, 38] We also observe a strong emission at 1.75 eV which has been attributed to radiative recombination of neutral excitons bound to defects (X^{o}) caused by Sulfur (S) vacancies. [39, 40]

7.2.3 Electrical Properties of MoS₂/Graphene

The conduction behavior of MoS₂ films deposited on Si/SiO₂ substrates was determined by fabricating field effect transistor (FET) devices with 50 nm thick Au as source and drain electrodes, 300 nm thick SiO₂ served as the dielectric layer, while doped silicon was used as the back gate. The output and transfer characteristics of our MoS₂ FET showed a *p*-type conduction behavior with a field effect mobility of $\approx 12.24 \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$ which is higher than mechanically exfoliated and CVD grown MoS₂ FETs on Si/SiO₂ substrates [41, 42] and I_{on/off} ratio of $\approx 10^6$.[31] The origin of p-type conduction in MoS₂/SiO₂ system is not explored in this research. However, Douli et al. in their studies showed that a p-type conduction in MoS₂ films on SiO₂ substrates could be due to the presence of localized trap states arising from impurities or defects within the oxide substrate or at the interface with the conducting channel, can redefine the effective Fermi level of the hybrid system to make it either n-type or p-type.[43] The trap states might have originated from the immobile ionic charges, SiO₂ surface oxygen dangling bonds or foreign impurities presented at the SiO₂/MoS₂ interface during MoS₂ synthesis. The defects present at the SiO₂ and MoS₂ interface could also be due to the 'Mo' diffusion into SiO₂ when it was sulfurized in the CVD furnace at the high temperature of 600 °C. [44, 45]. When a semi-metallic graphene is assembled over semiconducting MoS₂ a Schottky junction like behavior is expected. The schematic of carrier transport at MoS₂/graphene junction under forward and reverse bias is shown in Figure 7.4 (a).



Figure 7.4 (a) Energy level alignment for MoS_2 / graphene structure, band diagram at junction in equilibrium, forward bias and reverse bias with carrier flow

Figure 7.5 (a) shows the current-voltage (J-V) characteristics measured at $MoS_2/graphene$. As can be seen from the J-V plot and the |log (J)| vs. V plot shown in the inset, the junction shows a clear rectification as exponentially increasing current passes in the forward bias while there is strong resistance in the reverse bias which prevents flow of holes from graphene to MoS_2 . From the current-voltage and the Raman measurement results which shows an influx of holes in to graphene from MoS_2 at the junction, it can be concluded that a Schottky barrier junction (SBJ) similar to a metal-semiconductor junction is formed between graphene and MoS_2 . However, we observe that the current never reaches a complete saturation in reverse bias unlike an ideal metal-semiconductor Schottky junction. This is due to the bias dependent Fermi level observed in graphene which results in a variable Schottky barrier height (SBH) at different bias voltages.



Figure 7.5 (a) J-V curve showing rectification under reverse bias and exponentially increasing current in forward bias (inset: Absolute value of J-V plotted in logarithmic scale) (b) J-V at different temperatures showing an increase in J with increase in temperature (c) $1/C^2$ plotted against voltage (d) Reverse current from experimental and analytical model

We have used the model developed by Tongay *et al.* for the current observed in graphene – semiconductor junctions taking into account the change in graphene Fermi level with applied bias. [46] The current in the Schottky diode is given by

$$J(V) = A * T^2 \exp\left(-\frac{e\Phi^0_{SBH} + e\Delta\Phi_{SBH}(V)}{k_B T}\right) \left[\exp\left(\frac{eV}{k_B T}\right) - 1\right]$$
(1)

Where A^* is Richardson's constant, Φ^0_{SBH} is Schottky barrier height (SBH) at zero bias and $e\Delta\Phi_{SBH}(V)$ is the difference in Schottky barrier height due to change in the Fermi level of Graphene with applied bias and is given by

$$e\Delta\Phi_{SBH}(V) = -(\Delta E_F^G(V)) = V_F[\sqrt{\pi(n_0 - n_{induced})} - \sqrt{\pi n_0}] = -\frac{1}{2} V_F \sqrt{\frac{\pi\varepsilon_s\varepsilon_0 N_A(V_0 + V_R)}{2en_0}}$$
(2)

where N_A is acceptor concentration in the semiconductor, V_0 is the built-in voltage, V_R reverse bias voltage and n_0 the carrier density in graphene. The modified saturation current is

$$J_s(V) = A^* T^2 \exp(-\frac{e\Phi^0_{SBH} + e\Delta\Phi_{SBH}(V)}{k_B T})$$
(3)

Zero bias Schottky barrier height (Φ_{SBH}^0) was measured from the temperature dependent J-V measurements carried at four different temperatures and capacitance- voltage (C-V measurements were used to measure acceptor concentration N_A and built in voltage V_0 . The change in current (|J|) with temperature is shown in Figure 7.5 c which shows that the current increases with temperature in both bias directions. The value of Φ_{SBH}^0 can be obtained from the slope of Richardson's plot given by [47, 48]

$$\Phi^{0}_{SBH} = \frac{V_1}{n} - \frac{k}{q} \frac{\partial [\ln(\frac{I}{T^2})]}{\partial(\frac{1}{T})}$$
(4)

We measured the Φ_{SBH}^0 of MGJ at V₁= 0.3 V to be 139.20 meV where the ideality factor '*n*' is obtained from intercept of the linear region of the *J*-*V* plot and measured to be 1.86 while Richardson's constant A* is 0.70 x 10⁻⁶ A cm⁻² K⁻². The greater than unity value of the ideality factor can be attributed to enhanced image-force lowering across the MGJ and Schottky-barrier inhomogeneity.[46] Capacitance-Voltage measurements are used to obtain built – in voltage (V_0) and doping density values, from Schottky-Mott relationship. When square of inverse of capacitance ($1/C^2$ (cm⁴F⁻²)) plotted against applied voltage in reverse bias, ($1/C^2$) varies linearly with applied bias before saturating this can be extrapolated to intersect the abscissa to obtain built-in potential $V_0 = (\Phi_{MOS_2} - \Phi_{Gr})$ and the slope of the linear region gives the doping density (N_A) of MoS₂. From Figure 7.5 d we can see that the measured values hold the linearity in reverse bias. We measured V_0 and N_A values to be 280 mV and 6.64 x 10²¹ cm⁻³ (3.32 x 10¹⁵ cm⁻²) respectively. By considering the work function of graphene from literature to be 4.50 eV, the work function of *p*-type MoS₂ is found to be 4.78 eV.[49] Experimental and theoretical values of reverse current obtained from the analytical model developed shown in figure 7.5 b closely match each other.

7.2.4 MoS₂/Graphene Photodetector

One of the most interesting application of Schottky diodes is their use as photodetectors. $MoS_2/graphene$ hybrid structure would be an excellent material for photo-detection as both materials show high photosensitivity. Graphene being a gapless material exhibits sensitivity to light over a wide range of wavelengths from visible light to mid and far infrared. High electron mobility and high absorption would also contribute towards high responsivity in graphene. On the other hand, atomically thin MoS_2 has shown some interesting optical phenomenon such as 10^4 fold increase in photoluminescence when

compared to bulk MoS₂ and significant responsivity in visible light.[38] By Stacking both MoS₂ and graphene together to form Schottky junction, we would be able to harness these features of both materials to build high performing photodetectors. In general Schottky photodiodes operates in two modes: (1) Energy gap excitation where electron-hole pairs are generated by band to band excitation in the semiconductor when energy of photon is greater than the band gap of semiconductor $hv > E_g$ and (2) Internal photoemission which is observed when the incident light energy is greater than the Schottky barrier energy and smaller than the bandgap of semiconductor $(\Phi_{SBH} < hv < E_g)$ resulting in photo emission of carriers from metal surface close to the junction which are then driven towards the semiconductor over the Schottky barrier due to the electric field in the junction. This leads to a wide spectral range of detection based on the SBH values obtained. [50] Both schemes are explained in figure 7.6 (b). A Schottky barrier photodiode is conventionally operated in reverse bias where the increase in barrier results in very small reverse current and hence any increase in current by the photocurrent can be detected with a high sensitivity. This makes it easy to exactly measure the photocurrent generated. The expression for diode current is given by

$$J(V) = A * T^{2} [\exp(\frac{eV}{k_{B}T}) - 1] - J_{op}$$
(5)

where J_{op} is the photon induced current equal to $R * P_{op}$ where R be the responsivity of the photodiode and P_{op} is the power of incident light, the first part of the equation is the current equation in ideal Schottky diode. By replacing the diode current term with the modified Schottky diode current presented in equation (1), we can write the diode current equation as

$$J(V) = A * T^{2} \exp(-\frac{e\Phi^{0}_{SBH} + e\Delta\Phi_{SBH}(V)}{k_{B}T}) [\exp(\frac{eV}{k_{B}T}) - 1] - J_{op}$$
(6)



Figure 7.6. (a) MGPD under illumination (b) Illustration of photocurrent generation in a Schottky diode under illumination (c) J-V at dark and under 100 mW/cm² illumination (d) Photo responsivity of MGPD at different wavelengths (Spectral Response) (e) Photoconductivity of graphene.

In order to study the operation of the photodiode in both modes we have studied the performance of the device in a wide spectral range from 400 nm to 1500 nm. The spectral response 'R' of the device measured at 2 V reverse bias is shown in figure 4d. When the diode is operated in energy gap excitation mode where $hv > E_g^{MoS_2}$ a maximum photo responsivity of $R_{max} = 0.52 \text{ AW}^{-1}$ is observed at 590 nm. Beyond 670.00 nm, where the incident photon energy is smaller than $E_g^{MoS_2}$ the photo current is due to internal photo

emission in Graphene where photo generated holes move towards MoS₂. We find that the responsivity actually increases from 980 nm reaching a maximum value of 1.26 AW⁻¹ at 1440 nm before starting to fall. The obtained responsivity is not the highest but higher than many state of the art devices. This increase in responsivity follows the increase in photoconductivity observed in graphene which is shown in figure 7.6 (e). This kind of increased optical conductivity in graphene has been observed in intra band conductivity in earlier studies. [51]

One of the most important detector characteristics is the noise equivalent power (NEP) which represents the lowest input power giving a unit signal/noise ratio. In a Schottky junction photodetector, both Johnson noise which is due to thermal noise associated with the resistance and shot noise which is associated with the overall current contribute towards the overall noise. The Johnson noise current i_j is given by $i_j = \sqrt{\frac{4kBT}{R_{in}}}$

where k is the Boltzmann constant, T is the absolute temperature and R_{eq} is the equivalent resistance. The R_{eq} is obtained from the slope of I-V of MGPD at -2 V measured in dark. For R_{eq} of 5.83 M Ω and for normalized bandwidth (B) at room temperature the Johnson noise i_j is 1.68 pA. The shot noise is given by $i_s = \sqrt{2q(I_{ph} + I_d)B}$ where I_{ph} is photo current and I_d is dark current. The measured shot noise current at -2V is 11.2 pA at 1400 nm. The total noise current i_n is 11.3 pA and the NEP which is the ratio of i_n and the responsivity is calculated.[52, 53] We obtain a minimum NEP of 7.8 x 10⁻¹² W/ \sqrt{Hz} and corresponding detectivity of 4.2 x 10¹⁰ cm. \sqrt{Hz}/W at 1400 nm. The measured NEP of MGPD is less than typical Si photodetectors.

7.3 Conclusions

We have fabricated P:ZnO core-shell/n-ZnO radial p-n junction based photodetector which showed a record photoresponsivity of 60 A/W under UV illumination. High junction area characteristic of core-shell structures and low recombination losses owing to both the design and method of synthesis contributes to the high responsivity observed. We also stacked few layers of p-type MoS₂ over few layers of graphene to form a Schottky junction photodiode with wide bandwidth. Stable optical properties of MoS2 and high carrier mobility in graphene contribute to efficient charge separation at the junction. Apart from the impressive charge carrier properties, graphene also exhibits a wide bandwidth absorption which allows for the photodiode to operate in both energy gap excitation mode and internal photo emission mode, there by extending the detection range to higher wavelengths which is not reported in earlier MoS₂/graphene Schottky junctions. From the spectral response, we find that the fabricated device performs well over a wide range of wavelengths from visible region to mid infrared region with internal photo emission in graphene playing a major role in extending the range of the device. We have studied the junction characteristics and presented the underlying phenomenon of charge transfer in 2D MoS₂/graphene Schottky junction. CVD process is used in growing large area MoS₂ and graphene which along with the simplicity of fabrication can be easily integrated with metal oxide semiconductor (CMOS) technology.

7.4 References

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CHAPTER 8

Summary and Future Direction

8.1 Summary

In summary, this dissertation has detailed the development of Sonochemical method to synthesize ZnO nanorods, ZnO-ZnO core-shell nanorods, ZnO-ZnS core-shell nanorods, 2D ZnO nanowalls, 2D ZnO nanoflakes and their applications in sensing and energy harvesting. The developed method is 10 times faster than other chemical methods and synthesizes high quality single crystalline nanostructures at ambient conditions. For the first time, both n- type and p-type doping in ZnO via. Sonochemical method was reported using 'Al' and 'P' as dopants respectively. We have for the first time fabricated a radial pn junction photo diode in a core-shell geometry. The core-shell structure was synthesized without using any external surfactants but rather used only different molar ratios of precursors. A doping concentration of 6.7×10¹⁷ cm⁻³ was achieved for the P:ZnO nanoshells. The as fabricated homogeneous p-n junction shows a rectifying behavior with an ideality factor of 37 and a reverse saturation current of 200 µA at -2 V. The developed method is fast, inexpensive and performed at ambient conditions while resulting very small amount of impurities. This facilitates the use of flexible substrates for ZnO nanostructures, opening the door for flexible and wearable electronic and photonic devices based on ZnO. The photoresponse of the device was studied under UV illumination and the device recorded 60 A/W photoresponsivity which is the highest in ZnO p-n junction UV detectors. The high responsivity values are attributed to the increased junction area in core-shell geometry and reduced recombination losses at the interface.

We have also explored future 2D materials such as MoS₂ and graphene and succeeded in fabricating a high bandwidth photodetectors by stacking few layers of p-type MoS₂ over few layers of graphene to form Schottky junction photodiode. Stable optical properties of MoS₂ and high carrier mobility in graphene contribute to efficient charge separation at the junction. Apart from the impressive charge carrier properties, graphene also exhibits a wide bandwidth absorption which allows for the photodiode to operate in both energy gap excitation mode and internal photo emission mode, there by extending the detection range to higher wavelengths which is not reported in earlier MoS₂/graphene Schottky junctions. From the spectral response, we find that the fabricated device performs well over a wide range of wavelengths from visible region to mid infrared region with internal photo emission in graphene playing a major role in extending the range of the device. We have studied the junction characteristics and presented the underlying phenomenon of charge transfer in 2D MoS_2 /graphene Schottky junction. CVD process is used in growing large area MoS₂ and graphene which along with the simplicity of fabrication can be easily integrated with metal oxide semiconductor (CMOS) technology.

Further we have fabricated vertically standing 1D ZnO-NRs and lateral 2D ZnO-NFs on Au coated Si, using sonochemical technique to prepare label free, highly sensitive and selective electrochemical cortisol immunosensors. Both 1D ZnO-NRs and 2D ZnO-NFs exhibit selective electrochemical cortisol detection at 1pM 100 times better than ELISA. ZnO-NFs with sharp and distinct oxidation peaks indicate a much better surface controlled process when compared to ZnO-NRs. The sensitivity of ZnO-NFs is 7.74 μ A/M which is comparable to 11.86 μ A/M of ZnO-NRs coupled with ease of fabrication in case of 2D materials, we present ZnO nanoflakes as an alternate material for use as

immobilizing matrix in electrochemical biosensors. It is demonstrated in this work that ZnO-NSs based immunosensors are capable of detecting cortisol in human saliva sample and well-validated using ELISA values. The results of presented work suggest ZnO-NSs with high surface-to-volume ratio (NRs) and larger surface area (NFs) as viable materials for cortisol sensing in human saliva.

The role of ZnO nanorods as anode in organic photovoltaic cells was also explored. We have used ZnO nanorods grown on FTO as anodes in fabricating DSSCs. As fabricated device has shown an efficiency of 2.74%, and high fill factors. However the devices suffered from low short circuit currents due to high recombination losses and low stability as ZnO degraded rapidly in the acidic dye. The leakage of electrolyte has also been detrimental factor. From the gained experience of fabricating DSSCs we have setup a perovskite solar cell fabrication setup, in an environment isolated from the ambience and used ZnO nanorods grown on FTO as anodes. The perovskite formation and filling at different spin speeds was studied using SEM and XRD. The as fabricated devices suffered from high recombination losses due to thick seed layer, which led us to alternative geometry where ZnO nanoflakes grown over graphene/ PET was used as seed layer for ZnO nanorod growth. The as synthesized seed layer is less than 200 nm thick and single crystalline in nature which enhances the electron transport.

8.2 Future Direction

From the knowledge and experience gained over the course of this dissertation we propose two applications for Sonochemically synthesized ZnO nanostructure

8.2.3 Flexible and Transparent P:ZnO Nanoshell-/n-ZnO Nanorod Core- Shell LED

Ultra violet light emitting diodes have attracted tremendous attention in wide range of applications. ZnO with wide bandgap of 3.37 eV and high exciton binding energy of 60 meV is one of the most promising materials for fabricating UV LEDs. Encouraged by the successful p-type doping and high photoresponsivity of P:ZnO nanoshell/ n-ZnO nanorod core- shell structures under UV illumination we propose to further characterize and optimize the P:ZnO shell thickness and length for optimum electroluminescence. The proposed device is shown in figure 8.1. Where the P:ZnO nanoshell- n-ZnO nanorod core-shell structures are grown on graphene coated flexible substrate (PET)



Figure 8.1 P:ZnO nanoshell- n-ZnO nanorod core- shell structure grown on Au/graphene

8.2.2 Flexible Perovskite Solar Cells

With the insight obtained from our study of ZnO nanorod based perovskite solar cells, we propose fabricating of flexible highly efficient perovskite solar cells by growing ZnO- ZnS

core shell nanorods on ZnO nanoflake/ graphene/PET substrate. ZnO- ZnS core shell nanostructures have been successfully used in DSSCs solar cells by various groups to decrease the recombination losses as ZnS with larger band gap than ZnO isolated the injected electrons from being recombined with the holes. The proposed device is shown in figure 8.2



Figure 8.2 Perovskite solar cells based on ZnO-ZnS core-shell nanorods synthesized on ZnO nanoflakes/graphene/ PET as electron transport layer

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